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# Numerical analysis of history-dependent hemivariational inequalities and applications to viscoelastic contact problems with normal penetration



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# ABSTRACT

In this paper numerical approximation of history-dependent hemivariational inequalities with constraint is considered, and corresponding Céa's type inequality is derived for error estimate. For a viscoelastic contact problem with normal penetration, an optimal order error estimate is obtained for the linear element method. A numerical experiment for the contact problem is reported which provides numerical evidence of the convergence order predicted by the theoretical analysis.

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## 1. Introduction

Hemivariational inequalities are inequality problems characterized by nonconvexity and nonsmoothness, which may arise in engineering, economics, contact mechanics, etc. Techniques from nonsmooth analysis are applied to deal with difficulties caused by the nonsmoothness, which plays an important role in the mathematical theory of hemivariational inequality. In particular, properties of the generalized (Clarke) directional derivative and subdifferential of a locally Lipschitz function are frequently used; see [1–3].

Numerical methods for hemivariational inequalities have attracted much interest recently, in which the finite element approximation is widely used. An early comprehensive reference in this direction is [4]. In [5], an elliptic hemivariational inequality with convex constraint is considered, and two unconstrained problems can be regarded as special cases of a constrained problem. Finite element approximation is used to solve the problem, and the numerical solutions tend on subsequences to the solution of continuous problems. In [6], a variational–hemivariational inequality, which involves a convex functional and a nonconvex functional, is analyzed under appropriate assumptions. The existence and uniqueness of a solution are presented, and strong convergence for the finite element solution to the inequality is proved. Furthermore, an optimal first-order error estimate is derived for the linear element method. In [7], a general form of elliptic hemivariational inequality with or without convex constraint is considered, internal approximation of the discrete problem is studied with

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https://doi.org/10.1016/j.camwa.2018.12.038 0898-1221/© 2019 Elsevier Ltd. All rights reserved. convergence result proved and Céa's type inequalities derived. The theoretical results are applied to some particular contact problems, and optimal error estimates are derived using the linear element method. A further study of the paper [7] is given in [8], where external approximation of the inequality is analyzed and new techniques are developed to derive optimal order error estimate. In [9], both internal and external approximations for a general stationary variational–hemivariational inequality are studied. The numerical solutions are proved to converge strongly under minimal solution regularity property available from the well-posedness analysis of the problem. Recently, semipermeable media is studied in the framework of hemivariational inequalities in [10]. The model covers both isotropic, homogeneous semipermeable media and non-isotropic, heterogeneous semipermeable media. Interior and boundary semipermeable terms are considered, and external approximation of an abstract hemivariational inequality is illustrated. Convergence and optimal order error estimate of numerical solutions are shown.

History-dependent hemivariational inequalities are documented in [11–18]. A hemivariational inequality defined on a whole space is studied theoretically in [13]. In [12], numerical methods of temporally semi-discrete and fully discrete schemes are introduced to approximate a quasivariational inequality. In [17], a discrete operator is introduced to approximate the history-dependent operator with specific form, and a fully discrete scheme is used to solve a variational-hemivariational inequality numerically. In [18], an existence and uniqueness result of a general variational-hemivariational inequality is proved. Its internal approximation is presented in [19], and an optimal order error bound is derived. In this paper we further study numerical methods of history-dependent hemivariational inequalities, including its external approximation. In Section 2, a general hemivariational inequality is introduced. Numerical method is used to solve the hemivariational inequality in Section 3. In Section 4, the theoretical results are applied to a contact model with normal penetration and an optimal error estimate is derived for the linear element method. Finally, a numerical example is reported providing numerical evidence of the convergence order predicted by the theoretical analysis in Section 5.

### 2. A general history-dependent hemivariational inequality

In this section we consider a general history-dependent hemivariational inequality and discuss the existence and uniqueness of the solution.

Let  $X, X_j$  be normed spaces with norms  $\|\cdot\|_X$  and  $\|\cdot\|_{X_j}$ , let  $X^*$  and  $X_j^*$  be the topological dual of X and  $X_j$ . The duality pairings are denoted by  $\langle \cdot, \cdot \rangle_{X^* \times X}$  and  $\langle \cdot, \cdot \rangle_{X_j^* \times X_j}$ , respectively. When no confusion may arise, the notation  $\langle \cdot, \cdot \rangle_{X^* \times X}$  may be replaced by  $\langle \cdot, \cdot \rangle$  for simplicity. Let  $\mathbb{R}_+ = [0, +\infty)$  and K be a nonempty subset of X.  $C(\mathbb{R}_+; X)$  and  $C(\mathbb{R}_+; K)$  denote the spaces of continuous functions from  $\mathbb{R}_+$  to X and  $\mathbb{R}_+$  to K respectively. Let there be given  $A : X \to X^*, S : C(\mathbb{R}_+; X) \to C(\mathbb{R}_+; X^*)$ ,  $\gamma_j : X \to X_j$  and function  $j : X_j \to \mathbb{R}$  is locally Lipschitz. Then the general hemivariational inequality to be studied is the following.

**Problem 1.** Find a function  $u \in C(\mathbb{R}_+; K)$  such that  $\forall t \in \mathbb{R}_+$ ,

$$\langle Au(t), v - u(t) \rangle + \langle (Su)(t), v - u(t) \rangle + j^0(\gamma_j u(t); \gamma_j v - \gamma_j u(t)) \geq \langle f(t), v - u(t) \rangle \quad \forall v \in K.$$

$$(2.1)$$

The following assumptions are adopted in the analysis of Problem 1.

(2.2)
(7 2)
(2.5)
(2.4)
(2.5)
(2.6)
(2.7)

$$f \in \mathcal{C}(\mathbb{R}_+; X^*). \tag{2.8}$$

We explain the assumptions as follows. For the locally Lipschitz function  $j : X_j \rightarrow \mathbb{R}$ , we use  $j^0(w; z)$  to denote the generalized (Clarke) directional derivative of j at  $w \in X_j$  in the direction  $z \in X_j$ , defined by

$$j^{0}(w; z) = \limsup_{y \to w, \lambda \downarrow 0} \frac{j(y + \lambda z) - j(y)}{\lambda}$$

In applications to contact mechanics, the function j usually denotes an integral over a contact boundary and  $X_j$  represents the space of square integrable functions or vector-valued square integrable functions over the contact boundary.

The subdifferential (generalized gradient) of j at w, denoted as  $\partial j(w)$ , is a subset of the dual space  $X_i^*$  defined by

$$\partial j(w) = \{ \xi \in X_j^* \mid j^0(w; z) \ge \langle \xi, z \rangle_{X_j^* \times X_j} \, \forall z \in X_j \}.$$

By the definition of the subdifferential, the generalized directional derivative can be calculated by the relation  $j^0(w; z) = \max\{\langle \xi, z \rangle_{X_j^* \times X_j} | \xi \in \partial j(w)\}$ . More properties about generalized directional derivative and generalized gradient can be found in [1–3].

The operator  $S : C(\mathbb{R}_+; X) \to C(\mathbb{R}_+; X^*)$  is called a history-dependent operator if for any  $n \in \mathbb{N}$ , there exists a constant  $s_n > 0$  such that  $\forall t \in [0, n]$ ,

$$\|(\mathcal{S}u_1)(t) - (\mathcal{S}u_2)(t)\|_{X^*} \le s_n \int_0^t \|u_1(s) - u_2(s)\|_X ds \quad \forall \, u_1, \, u_2 \in C(\mathbb{R}_+; X).$$
(2.9)

History-dependence refers to the fact that the value (Su)(t) at a time t depends on the values of u in the entire interval [0, t]. The operator  $A : X \to X^*$  is called pseudomonotone if it is bounded and  $u_n \to u$  weakly in X together with  $\limsup \langle Au_n, u_n - u \rangle \leq 0$  imply

$$\langle Au, u - v \rangle < \liminf \langle Au_n, u_n - v \rangle \quad \forall v \in X$$

A is strongly monotone if there exists  $m_A > 0$  such that

$$\langle Av_1 - Av_2, v_1 - v_2 \rangle \ge m_A \|v_1 - v_2\|_X^2 \quad \forall v_1, v_2 \in X$$

The closed, convex set *K* reflects the constraint of the problem. For the contact problem we study, *K* usually denotes the admissible displacement field satisfying  $v_{\nu} \leq g$  on the contact boundary, where  $v_{\nu}$  is the normal component of the displacement  $\mathbf{v}$  and  $g \geq 0$  represents the upper bound of penetration.

Hypothesis (2.3)(a) means that the norm of the subdifferential of *j* is bounded by a linear function; (2.3)(b) describes constraint to the generalized directional derivative of *j* which guarantees the uniqueness of the solution to Problem 1. In addition, (2.3)(b) is equivalent to the relaxed monotone condition

$$\langle \partial j(z_1) - \partial j(z_2), z_1 - z_2 \rangle \ge -\alpha_j \|z_1 - z_2\|_{X_j}^2 \quad \forall z_1, z_2 \in X_j;$$

see [3,7,14,17]. Using (2.3),  $0 \in K$  and the strong monotonicity of A, we get that there exists a constant  $c \ge 0$  such that  $\langle Av, v \rangle \ge m_A \|v\|_X^2 - c\|v\|_X$  for all  $v \in X$ .

The following existence and uniqueness result for Problem 1 can be derived by taking  $\varphi((\mathcal{S}u)(t), u(t), v) = \langle (\mathcal{S}u)(t), v \rangle$ in [18, Theorem 5] and modifying slightly its proof.

**Theorem 2.** Assume that conditions (2.2)–(2.8) hold. Then Problem 1 has a unique solution  $u \in C(\mathbb{R}_+; K)$ .

## 3. Numerical method

In this section a fully discrete method is introduced and studied for the numerical solution of Problem 1. We use finite dimensional set  $K^h$  to approximate the convex subset K and  $K^h$  is not necessarily a subset of K. In the case  $K^h \subset K$ , the approximation is called internal, such a numerical method is studied in solving a variational–hemivariational inequality in [19]. In the case  $K^h \nsubseteq K$ , the corresponding approximation is called external. A Céa's type inequality is derived for external approximation of Problem 1.

Next we describe the discrete scheme for Problem 1, in which the spatial variable, the time variable and the historydependent operator are all discretized. On the spatial discretization, a regular family of finite element partition  $\{T^h\}$  is introduced where *h* is the mesh parameter. Corresponding finite element space  $X^h \subset X$  and  $K^h \subset X^h$  are presented, and  $K^h$ is assumed to be convex and closed with  $0 \in K^h$ . On the temporal discretization, *k* is the time step size and uniform partition is used for the interval I = [0, T]. Denote  $k = \frac{T}{N}$  and  $t_n = nk$  where  $0 \le n \le N$ .

Now we describe the discretization for history-dependent operator. Let  $\mathcal{L}(X)$  and  $\mathcal{L}(X; X^*)$  be the spaces of bounded linear operators from X to X and X to  $X^*$  respectively. Let  $C(\mathbb{R}_+ \times \mathbb{R}_+; \mathcal{L}(X))$  and  $C^2(\mathbb{R}_+ \times \mathbb{R}_+; \mathcal{L}(X))$  be the spaces of continuous and twice continuously differentiable functions defined on  $\mathbb{R}_+ \times \mathbb{R}_+$  with values in  $\mathcal{L}(X)$ . We focus on  $S : C(\mathbb{R}_+; X) \to C(\mathbb{R}_+; X^*)$  with a specific form

$$(\mathcal{S}v)(t) = \mathcal{R}\left(\int_0^t q(t,s)v(s)ds + a_s\right) \quad \forall v \in C(\mathbb{R}_+;X), \ \forall t \in \mathbb{R}_+,$$
(3.1)

where  $\mathcal{R} \in \mathcal{L}(X; X^*)$ ,  $q \in C(\mathbb{R}_+ \times \mathbb{R}_+; \mathcal{L}(X))$ ,  $a_S \in X$ . The above operator S is useful in dealing with history-dependent problems. For example, the well-known Volterra operator S, defined by

$$(Sv)(t) = \int_0^t \mathcal{B}(t-s)v(s)ds \quad \forall v \in C(\mathbb{R}_+;X), \ \forall t \in \mathbb{R}_+,$$

is a history-dependent operator, where  $\mathcal{B} \in C(0, T; \mathcal{L}(X; X^*))$ . More discussions on  $\mathcal{S}$  can be founded in [12,13,17,20]. For the integral in (3.1), we apply the trapezoidal rule

$$\int_0^{t_n} Z(s) ds \approx k \sum_{i=0}^n {}^{\prime} Z(t_i),$$

where a prime indicates that the coefficients of the first and last terms of the summation are to be halved. For n = 0, the summation is defined to be 0. Then the fully discrete operator  $S_n^{kh}$  is defined by

$$S_n^{kh} u^{kh} := \mathcal{R}\left(k \sum_{i=0}^{n'} q(t_n, t_i) u_i^{kh} + a_S\right), \ u^{kh} = \{u_i^{kh}\}_{i=0}^N \subset X^h.$$
(3.2)

Below, we denote  $||\mathcal{R}|| = ||\mathcal{R}||_{\mathcal{L}(X;X^*)}$  and  $||q|| = ||q||_{C(I \times I;\mathcal{L}(X))}$ . The space  $C^2(\mathbb{R}_+ \times \mathbb{R}_+; \mathcal{L}(X))$  is needed in bounding the difference between the history-dependent operator and its discrete analogue. In addition we introduce a constant *C* with possibly different values in different places which is independent of the mesh parameter *h* and the time step *k*. Now we illustrate the following fully discrete approximation of Problem 1.

**Problem 3.** The fully-discrete scheme for Problem 1 is to find the discrete solution  $u^{kh} := \{u_n^{kh}\}_{n=0}^N \subset K^h$  such that for  $0 \le n \le N$ ,

$$\begin{aligned} \langle Au_n^{kh}, v^h - u_n^{kh} \rangle &+ \langle \mathcal{S}_n^{kh} u^{kh}, v^h - u_n^{kh} \rangle + j^0 (\gamma_j u_n^{kh}; \gamma_j v^h - \gamma_j u_n^{kh}) \\ &\geq \langle f_n, v^h - u_n^{kh} \rangle \quad \forall v^h \in K^h. \end{aligned}$$

$$(3.3)$$

Under assumptions (2.2)–(2.3), (2.5)–(2.8), (3.1), 3.2 and a small time step

$$k < \frac{m_A - \alpha_j c_j^2}{\|\mathcal{R}\| \|q\|},\tag{3.4}$$

then there exists a unique solution  $u^{kh}$  of Problem 3. The proof is similar to that found in [19] where the solution existence and uniqueness of a discrete scheme is proved for a history-dependent variational-hemivariational inequality.

As a preparation for error estimation, we prove that the sequence  $||u_n^{kh}||_X$  is uniformly bounded.

**Proposition 4.** There exist constants M > 0 and  $k_0 > 0$  such that  $||u_n^{kh}||_X \le M \forall 0 < k \le k_0, \forall h > 0$ .

**Proof.** Taking  $v^h = 0$  in (3.3), we have

$$\langle Au_n^{kh}, u_n^{kh} \rangle \le \langle \mathcal{S}_n^{kh} u^{kh}, -u_n^{kh} \rangle + j^0 (\gamma_j u_n^{kh}; -\gamma_j u_n^{kh}) + \langle f_n, u_n^{kh} \rangle.$$

$$(3.5)$$

Note that  $0 \in K$ , thus for all  $v \in K$ ,

$$m_{A} \|v\|_{X}^{2} \le \langle Av - A0, v \rangle \le \langle Av, v \rangle + \|A0\|_{X^{*}} \|v\|_{X}.$$
(3.6)

Taking  $v = u_n^{kh}$ , we get

$$\langle Au_n^{kh}, u_n^{kh} \rangle \ge m_A \|u_n^{kh}\|_X^2 - \|A0\|_{X^*} \|u_n^{kh}\|_X.$$
(3.7)

From the formula 3.2, we have

$$\langle \mathcal{S}_{n}^{kh} u^{kh}, -u_{n}^{kh} \rangle \leq \| \mathcal{S}_{n}^{kh} u^{kh} \|_{X^{*}} \| u_{n}^{kh} \|_{X}$$

$$\leq k \| \mathcal{R} \| \| q \| \left( \sum_{i=0}^{n} \| u_{i}^{kh} \|_{X} \right) \| u_{n}^{kh} \|_{X} + \| \mathcal{R} \| \| a_{S} \|_{X} \| u_{n}^{kh} \|_{X}.$$

$$(3.8)$$

Take  $z_1 = \gamma_j u_n^{kh}$ ,  $z_2 = 0$  in (2.3)(b) and apply (2.3)(a) to obtain

$$j^{0}(\gamma_{j}u_{n}^{kh}; -\gamma_{j}u_{n}^{kh}) \leq \alpha_{j}\|\gamma_{j}u_{n}^{kh}\|_{X_{j}}^{2} - j^{0}(0; \gamma_{j}u_{n}^{kh}) \\ \leq \alpha_{j}c_{j}^{2}\|u_{n}^{kh}\|_{X}^{2} + c_{0}c_{j}\|u_{n}^{kh}\|_{X}.$$

$$(3.9)$$

Obviously,

$$\langle f_n, u_n^{kh} \rangle \le \|f_n\|_{X^*} \|u_n^{kh}\|_X. \tag{3.10}$$

Combining (3.5) and (3.7)-(3.10), we have

$$\begin{aligned} \|u_n^{kh}\|_X &\leq \frac{1}{m_A - \alpha_j c_j^2} (\|\mathcal{R}\| \|a_S\|_X + c_0 c_j + \|A0\|_{X^*} + \|f_n\|_{X^*}) \\ &+ \frac{k}{m_A - \alpha_j c_j^2} \|\mathcal{R}\| \|q\| \sum_{i=0}^n \|u_i^{kh}\|_X. \end{aligned}$$
(3.11)

We recall the discrete Gronwall's inequality [21]: let  $\{e_n\}_{n=0}^N$  and  $\{g_n\}_{n=0}^N$  be two sequences of non-negative numbers with

$$e_n \leq C_1 g_n + C_2 k \sum_{i=0}^n e_i, \quad 0 \leq n \leq N$$

where  $C_1$ ,  $C_2$  are two constants. Then there is a constant  $C_3$  such that for k > 0 sufficiently small, we have

$$\max_{0\leq n\leq N}e_n\leq C_3\max_{0\leq n\leq N}g_n.$$

Thus we obtain

$$\max_{0 \le n \le N} \|u_n^{kh}\|_X \le C \left( \max_{0 \le n \le N} \|f_n\|_{X^*} + \|\mathcal{R}\| \|a_S\|_X + c_0 c_j + \|A0\|_{X^*} \right).$$
(3.12)

Since  $f \in C(\mathbb{R}_+; X^*)$ , the sequence  $\{u_n^{kh}\}$  is bounded uniformly.  $\Box$ 

To proceed further for error analysis, we assume the operator  $A : X \to X^*$  is Lipschitz continuous, i.e., for a constant  $L_A > 0$ ,

$$\|Au - Av\|_{X^*} \le L_A \|u - v\|_X \quad \forall u, v \in X.$$
(3.13)

Let  $t = t_n$  in (2.1), we have

$$\langle Au_n, v - u_n \rangle + \langle S_n u, v - u_n \rangle + j^0(\gamma_j u_n; \gamma_j v - \gamma_j u_n) \ge \langle f_n, v - u_n \rangle.$$
(3.14)

where  $S_n u = \mathcal{R}\left(\int_0^{t_n} q(t_n, s)u(s)ds + a_S\right)$  and  $u_n = u(t_n)$ . Define a residual term

$$E_n(v) = \langle Au_n, v \rangle + \langle S_n u, v \rangle + j^0(\gamma_j u_n; \gamma_j v) - \langle f_n, v \rangle,$$
(3.15)

which play an important role in error estimation. Now we derive a Céa's type inequality.

**Theorem 5.** Assume (2.2)–(2.3), (2.5)–(2.8), (3.1), 3.2, (3.4) and (3.13). Furthermore, assume regularity conditions  $q \in C^2(\mathbb{R}_+ \times \mathbb{R}_+; \mathcal{L}(X))$  and  $u \in W^{2,\infty}_{loc}(\mathbb{R}_+; X)$ . Then the following error estimate holds for the numerical solution of Problem 1 defined by Problem 3:

$$\max_{0 \le n \le N} \|u_n - u_n^{kh}\|_X \le C \max_{0 \le n \le N} \left\{ \inf_{v^h \in K^h} \left( \|u_n - v^h\|_X + \|\gamma_j u_n - \gamma_j v^h\|_{X_j}^{\frac{1}{2}} + |E_n(v^h - u_n)|^{\frac{1}{2}} \right) + \inf_{v \in K} |E_n(v - u_n^{kh})|^{\frac{1}{2}} \right\}$$

$$+ Ck^2 \|u\|_{W^{2,\infty}(I;X)}.$$
(3.16)

Proof. We begin with

$$m_{A} \|u_{n} - u_{n}^{kh}\|_{X}^{2} \leq \langle Au_{n} - Au_{n}^{kh}, u_{n} - u_{n}^{kh} \rangle$$

$$= \langle Au_{n} - Au_{n}^{kh}, u_{n} - v^{h} \rangle + \langle Au_{n}, v^{h} - u_{n} \rangle + \langle Au_{n}, u_{n} - v \rangle$$

$$+ \langle Au_{n}, v - u_{n}^{kh} \rangle + \langle Au_{n}^{kh}, u_{n}^{kh} - v^{h} \rangle.$$
(3.17)

Combining (3.17), (3.3) and (3.14), we have

$$\begin{split} m_{A} \|u_{n} - u_{n}^{kh}\|_{X}^{2} &\leq \langle Au_{n} - Au_{n}^{kh}, u_{n} - v^{h} \rangle + \langle Au_{n}, v^{h} - u_{n} \rangle + \langle Au_{n}, v - u_{n}^{kh} \rangle \\ &+ \langle S_{n}u, v - u_{n} \rangle + j^{0}(\gamma_{j}u_{n}; \gamma_{j}v - \gamma_{j}u_{n}) - \langle f_{n}, v - u_{n} \rangle \\ &+ \langle S_{n}^{kh}u^{kh}, v^{h} - u_{n}^{kh} \rangle + j^{0}(\gamma_{j}u_{n}^{kh}; \gamma_{j}v^{h} - \gamma_{j}u_{n}^{kh}) \\ &- \langle f_{n}, v^{h} - u_{n}^{kh} \rangle \\ &= \langle Au_{n} - Au_{n}^{kh}, u_{n} - v^{h} \rangle + E_{n}(v^{h} - u_{n}) + E_{n}(v - u_{n}^{kh}) \\ &+ E_{S}(v, v^{h}, u_{n}, u_{n}^{kh}) + E_{i}(v, v^{h}, u_{n}, u_{n}^{kh}), \end{split}$$
(3.18)

where  $E_n(v^h - u_n)$  and  $E_n(v - u_n^{kh})$  are defined by (3.15),  $E_S(v, v^h, u_n, u_n^{kh})$  and  $E_j(v, v^h, u_n, u_n^{kh})$  are defined by

$$E_{S}(v, v^{h}, u_{n}, u_{n}^{kh}) = \langle S_{n}u - S_{n}^{kh}u^{kh}, u_{n}^{kh} - v^{h} \rangle$$
  
=  $\langle S_{n}u - S_{n}^{kh}u^{kh}, u_{n}^{kh} - u_{n} \rangle + \langle S_{n}u - S_{n}^{kh}u^{kh}, u_{n} - v^{h} \rangle,$  (3.19)

and

$$E_{j}(v, v^{h}, u_{n}, u_{n}^{kh}) = j^{0}(\gamma_{j}u_{n}; \gamma_{j}v - \gamma_{j}u_{n}) + j^{0}(\gamma_{j}u_{n}^{kh}; \gamma_{j}v^{h} - \gamma_{j}u_{n}^{kh}) - j^{0}(\gamma_{j}u_{n}; \gamma_{j}v^{h} - \gamma_{j}u_{n}) - j^{0}(\gamma_{j}u_{n}; \gamma_{j}v - \gamma_{j}u_{n}^{kh}).$$
(3.20)

Next we bound  $E_S(v, v^h, u_n, u_n^{kh})$  and  $E_j(v, v^h, u_n, u_n^{kh})$ . For any small  $\varepsilon > 0$ , we have

$$E_{S} \leq \|S_{n}u - S_{n}^{kh}u^{kh}\|_{X^{*}}\|u_{n} - u_{n}^{kh}\|_{X} + \|S_{n}u - S_{n}^{kh}u^{kh}\|_{X^{*}}\|u_{n} - v^{h}\|_{X}$$
  
$$\leq \varepsilon \|u_{n} - u_{n}^{kh}\|_{X}^{2} + \|u_{n} - v^{h}\|_{X}^{2} + \frac{1}{4}(1 + \frac{1}{\varepsilon})\|S_{n}u - S_{n}^{kh}u^{kh}\|_{X^{*}}^{2}.$$
(3.21)

Using the subadditivity of the generalized directional derivative, we get

$$\begin{split} E_{j} &\leq j^{0}(\gamma_{j}u_{n};\gamma_{j}v-\gamma_{j}u_{n}^{kh})+j^{0}(\gamma_{j}u_{n};\gamma_{j}u_{n}^{kh}-\gamma_{j}u_{n})+j^{0}(\gamma_{j}u_{n}^{kh};\gamma_{j}v^{h}-\gamma_{j}u_{n})\\ &+j^{0}(\gamma_{j}u_{n}^{kh};\gamma_{j}u_{n}-\gamma_{j}u_{n}^{kh})-j^{0}(\gamma_{j}u_{n};\gamma_{j}v^{h}-\gamma_{j}u_{n})-j^{0}(\gamma_{j}u_{n};\gamma_{j}v-\gamma_{j}u_{n}^{kh})\\ &\leq \alpha_{j}\|\gamma_{j}u_{n}-\gamma_{j}u_{n}^{kh}\|_{X_{j}}^{2}+j^{0}(\gamma_{j}u_{n}^{kh};\gamma_{j}v^{h}-\gamma_{j}u_{n})-j^{0}(\gamma_{j}u_{n};\gamma_{j}v^{h}-\gamma_{j}u_{n})\\ &\leq \alpha_{j}c_{j}^{2}\|u_{n}-u_{n}^{kh}\|_{X}^{2}+(2c_{0}+c_{1}\|\gamma_{j}u_{n}^{kh}\|_{X_{j}}+c_{1}\|\gamma_{j}u_{n}\|_{X_{j}})\|\gamma_{j}u_{n}-\gamma_{j}v^{h}\|_{X_{j}}. \end{split}$$

Since A is Lipschitz continuous, we have

$$\langle Au_{n} - Au_{n}^{kh}, u_{n} - v^{h} \rangle \leq L_{A} \|u_{n} - u_{n}^{kh}\|_{X} \|u_{n} - v^{h}\|_{X}$$

$$\leq \varepsilon \|u_{n} - u_{n}^{kh}\|_{X}^{2} + \frac{L_{A}^{2}}{4\varepsilon} \|u_{n} - v^{h}\|_{X}^{2}.$$

$$(3.22)$$

Thus, from (3.18), we obtain

$$\begin{split} m_{A} \|u_{n} - u_{n}^{kh}\|_{X}^{2} &\leq (2\varepsilon + \alpha_{j}c_{j}^{2})\|u_{n} - u_{n}^{kh}\|_{X}^{2} + (1 + \frac{L_{A}^{2}}{4\varepsilon})\|u_{n} - v^{h}\|_{X}^{2} \\ &+ (2c_{0} + c_{1}c_{j}\|u_{n}^{kh}\|_{X} + c_{1}c_{j}\|u_{n}\|_{X})\|\gamma_{j}u_{n} - \gamma_{j}v^{h}\|_{X_{j}} \\ &+ |E_{n}(v^{h} - u_{n})| + |E_{n}(v - u_{n}^{kh})| + \frac{1}{4}(1 + \frac{1}{\varepsilon})\|\mathcal{S}_{n}u - \mathcal{S}_{n}^{kh}u^{kh}\|_{X^{*}}^{2}. \end{split}$$

Using the conditions  $u \in C(I; K)$  and (2.7) and applying Proposition 4, we find

$$\begin{aligned} \|u_{n} - u_{n}^{kh}\|_{X} &\leq C \left( \|u_{n} - v^{h}\|_{X} + \|\gamma_{j}u_{n} - \gamma_{j}v^{h}\|_{X_{j}}^{\frac{1}{2}} + |E_{n}(v^{h} - u_{n})|^{\frac{1}{2}} \\ &+ |E_{n}(v - u_{n}^{kh})|^{\frac{1}{2}} + \|S_{n}u - S_{n}^{kh}u^{kh}\|_{X^{*}} \right). \end{aligned}$$

$$(3.23)$$

Recall the assumptions  $u \in W^{2,\infty}_{loc}(\mathbb{R}_+; X)$ ,  $q \in C^2(\mathbb{R}_+ \times \mathbb{R}_+; \mathcal{L}(X))$ . We have

$$\begin{split} \|\mathcal{S}_{n}^{kh}u^{kh} - \mathcal{S}_{n}u\|_{X^{*}} &\leq \|\mathcal{R}(k\sum_{i=0}^{n}{}^{\prime}q(t_{n},t_{i})u_{i}^{kh} + a_{S}) - \mathcal{R}(k\sum_{i=0}^{n}{}^{\prime}q(t_{n},t_{i})u_{i} + a_{S})\|_{X^{*}} \\ &+ \|\mathcal{R}(k\sum_{i=0}^{n}{}^{\prime}q(t_{n},t_{i})u_{i} + a_{S}) - \mathcal{R}(\int_{0}^{t_{n}}q(t_{n},s)u(s)ds + a_{S})\|_{X^{*}} \\ &\leq k\|\mathcal{R}\|\|q\|\left(\sum_{i=0}^{n}\|u_{i}^{kh} - u_{i}\|_{X}\right) + Ck^{2}\|u\|_{W^{2,\infty}(l;X)}. \end{split}$$

Thus, from (3.23),

$$\begin{split} \|u_n - u_n^{kh}\|_X &\leq C \left( \|u_n - v^h\|_X + \|\gamma_j u_n - \gamma_j v^h\|_{X_j}^{\frac{1}{2}} + |E_n(v^h - u_n)|^{\frac{1}{2}} \\ &+ |E_n(v - u_n^{kh})|^{\frac{1}{2}} + k^2 \|u\|_{W^{2,\infty}(I;X)} \right) + Ck \sum_{i=0}^n \|u_i - u_i^{kh}\|_X. \end{split}$$

Using discrete Gronwall's inequality again, we obtain the error bound (3.17).  $\Box$ 

In the case  $K^h \subset K \cap X^h$ , the approximation is internal and the residual term  $\inf_{v \in K} |E_n(v - u_n^{kh})| = 0$ . Then Theorem 5 reduces to the following result.

**Theorem 6.** Assume (2.2)–(2.3), (2.5)–(2.8), (3.1), 3.2, (3.4) and (3.13). Furthermore, assume  $K^h \subset K \cap X^h$  and regularities  $q \in C^2(\mathbb{R}_+ \times \mathbb{R}_+; \mathcal{L}(X))$  and  $u \in W^{2,\infty}_{loc}(\mathbb{R}_+; X)$ . Then we have the error estimate of the numerical solution of Problem 1 defined

by Problem 3:

$$\max_{0 \le n \le N} \|u_n - u_n^{kh}\|_X \le C \max_{0 \le n \le N} \left\{ \inf_{v^h \in K^h} \left( \|u_n - v^h\|_X + \|\gamma_j u_n - \gamma_j v^h\|_{X_j}^{\frac{1}{2}} + |E_n(v^h - u_n)|^{\frac{1}{2}} \right) \right\} + Ck^2 \|u\|_{W^{2,\infty}(I;X)}.$$
(3.24)

A Céa's type inequality of form (3.24) is derived in [19] for the error of approximations of a history-dependent variationalhemivariational inequality.

#### 4. Application to a viscoelastic contact problem

In this section a representative viscoelastic contact model with normal penetration is considered. The numerical method and theoretical results illustrated in Section 3 are applied to the contact problem, and an optimal order error estimate is derived for the linear finite element method.

We start with some standard notation in mechanics. Let d = 2or3 be the spatial dimension. Let  $\Omega \subset \mathbb{R}^d$  be a Lipschitz domain, representing the reference configuration of the deformable body. Let  $\mathbf{u} = (u_i)$  be the displacement field, and  $\sigma = (\sigma_{ij})$  the stress tensor of second-order. They are the main unknowns of the problem. Symbols  $\mathcal{A} = (a_{ijkl})$  and  $\mathcal{B} = (b_{ijkl})$  both represent fourth-order tensors. All the subscripts i, j, k, l satisfy  $1 \leq i, j, k, l \leq d$ . Since the boundary  $\Gamma$  of  $\Omega$  is Lipschitz continuous, the unit outward normal  $\mathbf{v} = (v_i)$  exists a.e. on  $\Gamma$ . For a vector field  $\mathbf{v}$  defined on  $\Gamma$ ,  $v_v := \mathbf{v} \cdot \mathbf{v}$  denotes its normal component and  $\mathbf{v}_\tau := \mathbf{v} - v_v \mathbf{v}$  its tangential component. Similarly for a stress field  $\sigma, \sigma_v := (\sigma \mathbf{v}) \cdot \mathbf{v}$  represents its normal component and  $\sigma_\tau := \sigma \mathbf{v} - \sigma_v \mathbf{v}$  represents its tangential component. Let  $\varepsilon(\mathbf{u}) = (\nabla \mathbf{u} + (\nabla \mathbf{u})^T)/2$  be the linearized strain tensor. The classical formulation of the viscoelastic contact problem is as follows.

**Problem 7.** Find a displacement  $\boldsymbol{u} : \Omega \times \mathbb{R}_+ \to \mathbb{R}^d$  and a stress field  $\boldsymbol{\sigma} : \Omega \times \mathbb{R}_+ \to \mathbb{S}^d$  such that

$$\boldsymbol{\sigma}(t) = A\boldsymbol{\varepsilon}(\boldsymbol{u}(t)) + \int_{0}^{t} \mathcal{B}(t-s)\boldsymbol{\varepsilon}(\boldsymbol{u}(s))ds \text{ in } \Omega \times \mathbb{R}_{+},$$

$$\tag{4.1}$$

$$\operatorname{Div} \boldsymbol{\sigma}(t) + \boldsymbol{f}_0(t) = \boldsymbol{0} \text{ in } \boldsymbol{\Omega} \times \mathbb{R}_+, \tag{4.2}$$

$$\mathbf{u}(t) = \mathbf{0} \text{ on } \Gamma_1 \times \mathbb{R}_+, \tag{4.3}$$

$$\boldsymbol{\sigma}(t)\boldsymbol{\nu} = \boldsymbol{f}_2(t) \text{ on } \boldsymbol{\Gamma}_2 \times \mathbb{R}_+, \tag{4.4}$$

$$\begin{aligned} u_{\nu}(t) &\leq g, \ \sigma_{\nu}(t) + \xi_{\nu}(t) \leq 0, \\ (\sigma_{\nu}(t) + \xi_{\nu}(t))(u_{\nu}(t) - g) &= 0, \end{aligned} \right\} \ \text{on} \ \Gamma_{3} \times \mathbb{R}_{+},$$
 (4.5)

$$\sigma_{\tau}(t) = \mathbf{0} \text{ on } \Gamma_3 \times \mathbb{R}_+.$$

$$\tag{4.6}$$

The model is illustrated graphically in Fig. 1 of Section 5. Note that a slight different model is studied in [18]. In the following we give the description of the physical setting. A layer of elastic material is laid on the rigid foundation, and a viscoelastic body which occupies  $\Omega$  is in contact with the foundation. The boundary  $\Gamma$  of the body is divided into three measurable parts  $\Gamma_1$ ,  $\Gamma_2$  and  $\Gamma_3$ , with meas( $\Gamma_1$ ) > 0. The whole contact process is quasi-static.

Eq. (4.1) is a constitutive law for viscoelastic material with long memory (cf. [14,17,18,22]), where A is a nonlinear elasticity tensor and B is a relaxation tensor. Relation (4.2) is the equilibrium equation, where Div is the divergence operator, and  $f_0$  is the density of volume forces acting on the body  $\Omega$ . Relations (4.3) and (4.4) mean that the body is fixed on  $\Gamma_1$  and the density of surface forces  $f_2$  acts on  $\Gamma_2$ . Eq. (4.5) describes the state of the contact process on  $\Gamma_3$ . Since the penetration is restricted by  $v_{\nu} \leq g$ , the set of admissible displacements is a convex and closed subset of  $H^1(\Omega; \mathbb{R}^d)$ . The normal stress is described by a Clarke subdifferential of a function  $j_{\nu}$ . Relation (4.6) means the contact process is frictionless on the boundary  $\Gamma_3$ .

To study Problem 7, we introduce the space  $V = \{ \mathbf{v} = (v_i) \in H^1(\Omega; \mathbb{R}^d) \mid \mathbf{v} = 0 \text{ a.e. on } \Gamma_1 \}$  for the displacement field, and the space  $\mathcal{H} = \{ \mathbf{\tau} = (\tau_{ij}) \in L^2(\Omega; \mathbb{S}^d) \mid \tau_{ij} = \tau_{ji}, 1 \le i, j \le d \}$  for the stress field, where  $\mathbb{S}^d$  denotes the space of symmetric tensor of second-order. Denote  $\mathbb{Q}_{\infty} = \{ \mathcal{E} = (\mathcal{E}_{ijkl}) \mid \mathcal{E}_{ijkl} = \mathcal{E}_{jikl} = \mathcal{E}_{klij} \in L^{\infty}(\Omega), 1 \le i, j, k, l \le d \}$  with the associated norm defined by  $\|\mathcal{E}\|_{\mathbb{Q}_{\infty}} = \sum_{1 \le i, k, l \le d} \|\mathcal{E}_{ijkl}\|_{L^{\infty}(\Omega)}$ . The canonical inner products in the spaces  $\mathbb{R}^d$ ,  $\mathbb{S}^d$  and  $\mathcal{H}$  are defined by  $(\mathbf{u}, \mathbf{v})_{\mathbb{R}^d} = \sum_{i=1}^d u_i v_i$  for all  $\mathbf{u}, \mathbf{v} \in \mathbb{R}^d$ ,  $(\sigma, \tau)_{\mathbb{S}^d} = \sum_{1 \le i, j \le d} \sigma_{ij} \tau_{ij}$  for all  $\sigma, \tau \in \mathbb{S}^d$  and  $(\sigma, \tau)_{\mathcal{H}} = \int_{\Omega} \sum_{1 \le i, j \le d} \sigma_{ij}(x) \tau_{ij}(x) dx$  for all  $\sigma, \tau \in \mathcal{H}$  respectively. Since meas $(\Gamma_1) > 0$ , V is a Hilbert space with the inner product  $(\mathbf{u}, \mathbf{v})_V = (\boldsymbol{\varepsilon}(\mathbf{u}), \boldsymbol{\varepsilon}(\mathbf{v}))_{\mathcal{H}}$  for all  $\mathbf{u}, \mathbf{v} \in V$ .

For any  $\mathbf{v} \in V$ , the trace of  $\mathbf{v}$  on the boundary  $\Gamma$  is still denoted by  $\mathbf{v}$ . The normal trace is  $v_{\nu} = \mathbf{v} \cdot \mathbf{v}$  and tangential trace  $v_{\tau} = \mathbf{v} - v_{\nu} \mathbf{v}$ .

Let  $K = \{ \mathbf{v} \in V | v_{\nu} \le g \text{ a.e. on } \Gamma_3 \}$ , where the thickness function g is nonnegative. It is clear that K is convex and closed in V and  $0 \in K$ .

The function  $j_{\nu} : \Gamma_3 \times \mathbb{R} \to \mathbb{R}$  satisfies the following assumptions:  $j_{\nu}(\mathbf{x}, \cdot)$  is locally Lipschitz on  $\mathbb{R}$  for a.e.  $\mathbf{x} \in \Gamma_3$ ;  $j_{\nu}(\cdot, r)$  is measurable on  $\Gamma_3$  for all  $r \in \mathbb{R}$ , and there exists  $\overline{e} \in L^2(\Gamma_3)$  such that  $j_{\nu}(\cdot, \overline{e}(\cdot)) \in L^1(\Gamma_3)$ ; there exist constants  $\overline{c}_0, \overline{c}_1 > 0$  such

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that  $|\partial j_{\nu}(\mathbf{x}, r)| \leq \overline{c}_0 + \overline{c}_1 |r|$  for a.e.  $\mathbf{x} \in \Gamma_3$ , for all  $r \in \mathbb{R}$ ; there exists a constant  $\overline{\alpha}_{\nu} \geq 0$  such that the sum of  $j_{\nu}^0(\mathbf{x}, r_1; r_2 - r_1)$  and  $j_{\nu}^0(\mathbf{x}, r_2; r_1 - r_2)$  is less than or equal to the quantity  $\overline{\alpha}_{\nu} |r_1 - r_2|^2$  for all  $r_1, r_2 \in \mathbb{R}$  for a.e.  $\mathbf{x} \in \Gamma_3$ . For the elasticity tensor  $\mathcal{A} : \Omega \times \mathbb{S}^d \to \mathbb{S}^d$ , we assume  $\mathcal{A}(\cdot, \boldsymbol{\varepsilon})$  is measurable on  $\Omega, \mathcal{A}(\mathbf{x}, \mathbf{0}) = \mathbf{0}, \mathcal{A}(\mathbf{x}, \cdot)$  is strongly

For the elasticity tensor  $\mathcal{A} : \Omega \times \mathbb{S}^d \to \mathbb{S}^d$ , we assume  $\mathcal{A}(\cdot, \boldsymbol{\varepsilon})$  is measurable on  $\Omega$ ,  $\mathcal{A}(\boldsymbol{x}, \boldsymbol{0}) = \boldsymbol{0}$ ,  $\mathcal{A}(\boldsymbol{x}, \cdot)$  is strongly monotone and Lipschitz continuous. We denote by  $m_{\mathcal{A}} > 0$  the constant in the strong monotonicity, and by  $L_{\mathcal{A}} > 0$  the Lipschitz constant. The fourth-order tensor  $\mathcal{B}$  satisfies  $\mathcal{B} \in C(\mathbb{R}_+; \mathbb{Q}_\infty)$ .

The density of body force  $f_0$  and the density of surface traction  $f_2$  satisfy  $f_0 \in C(\mathbb{R}_+; L^2(\Omega; \mathbb{R}^d))$  and  $f_2 \in C(\mathbb{R}_+; L^2(\Gamma_2; \mathbb{R}^d))$  respectively. Using Riesz representation theorem, we can define a function  $f : \mathbb{R}_+ \to V^*$  by

$$\langle \boldsymbol{f}(t), \boldsymbol{v} \rangle_{V^* \times V} = (\boldsymbol{f}_0(t), \boldsymbol{v})_{L^2(\Omega; \mathbb{R}^d)} + (\boldsymbol{f}_2(t), \boldsymbol{v})_{L^2(\Gamma_2; \mathbb{R}^d)} \quad \forall \ \boldsymbol{v} \in V, \ t \in \mathbb{R}_+.$$

Then the weak formulation of Problem 7 is the following.

**Problem 8.** Find a displacement  $\boldsymbol{u} : \mathbb{R}_+ \to K$  such that  $\forall t \in \mathbb{R}_+$ ,

$$\left( \mathcal{A}\boldsymbol{\varepsilon}(\boldsymbol{u}(t)), \boldsymbol{\varepsilon}(\boldsymbol{v}) - \boldsymbol{\varepsilon}(\boldsymbol{u}(t)) \right)_{\mathcal{H}} + \left( \int_{0}^{t} \mathcal{B}(t-s)\boldsymbol{\varepsilon}(\boldsymbol{u}(s)) ds, \boldsymbol{\varepsilon}(\boldsymbol{v}) - \boldsymbol{\varepsilon}(\boldsymbol{u}(t)) \right)_{\mathcal{H}} + \int_{\Gamma_{3}} j_{\upsilon}^{0}(u_{\upsilon}(t); \upsilon_{\upsilon} - u_{\upsilon}(t)) d\Gamma \geq \langle \boldsymbol{f}(t), \boldsymbol{v} - \boldsymbol{u}(t) \rangle_{V^{*} \times V} \, \forall \, \boldsymbol{v} \in K.$$

$$(4.7)$$

We define the operators  $A: V \to V^*$ ,  $S: C(\mathbb{R}_+; V) \to C(\mathbb{R}_+; V^*)$  and function  $j: X_j \to \mathbb{R}$  as follows:

$$\langle A\boldsymbol{u},\boldsymbol{v}\rangle_{V^*\times V} = \left(\mathcal{A}\boldsymbol{\varepsilon}(\boldsymbol{u}),\boldsymbol{\varepsilon}(\boldsymbol{v})\right)_{\mathcal{H}} \quad \forall \,\boldsymbol{u},\boldsymbol{v}\in V,$$

$$(4.8)$$

$$(\mathcal{S}\boldsymbol{u})(t) = \int_0^t \mathcal{B}(t-s)\boldsymbol{\varepsilon}(\boldsymbol{u}(s))ds \quad \forall \, \boldsymbol{u} \in C(\mathbb{R}_+; V).$$
(4.9)

Let X = V,  $X_j = L^2(\Gamma_3)$ ,  $\gamma_j \boldsymbol{v} = v_v \forall \boldsymbol{v} \in V$ . Define

$$j(v) = \int_{\Gamma_3} j_v(v) d\Gamma \quad \forall v \in X_j.$$
(4.10)

Thus we have

$$j(\gamma_j \boldsymbol{\nu}) = \int_{\Gamma_3} j_{\boldsymbol{\nu}}(\boldsymbol{\nu}_{\boldsymbol{\nu}}) d\Gamma \quad \forall \, \boldsymbol{\nu} \in V.$$
(4.11)

Using the argument in [18, Section 6] and Theorem 2, under the hypotheses for Problem 7 and the smallness condition  $\overline{\alpha}_{\nu} \|\gamma_{j}\|^{2} < m_{\mathcal{A}}$ , we can show that Problem 8 has a unique solution  $\boldsymbol{u} \in C(\mathbb{R}_{+}; K)$ .

Next we consider a fully discrete method to solve Problem 8. For simplicity, assume  $\Omega$  is a polygonal/polyhedral. Let  $\{\mathcal{T}^h\}$  be a regular family of partition of  $\Omega$  into triangles/tetrahedrons, and the boundary  $\Gamma$  is divided into  $\Gamma_{i,l}$ ,  $1 \le l \le l_i$ ,  $1 \le i \le 3$ . If there is intersection between the side/face of an element and the set  $\Gamma_{i,l}$ , and the measurement of the intersection is positive with respect to  $\Gamma_{i,l}$ , then the side/face locates entirely in  $\Gamma_{i,l}$ . Denote  $V^h = \{\mathbf{v}^h \in C(\overline{\Omega})^d \mid \mathbf{v}^h|_T \in \mathbb{P}_1(T)^d$  for  $T \in \mathcal{T}^h, \mathbf{v}^h = \mathbf{0}$  on  $\Gamma_1$  be the linear element spaces for  $\{\mathcal{T}^h\}$ , and  $K^h$  be the subsets of  $V^h$  defined by  $K^h = \{\mathbf{v}^h \in V^h \mid v_v^h \le g$  at node points on  $\Gamma_3$ . Thus  $0 \in K^h$  is satisfied automatically. If the thickness function g is concave, then  $K^h \subset K$ . The numerical method for this case is studied in [19]. We consider the general case  $K^h \notin K$  in this paper.

For the time interval [0, T] the uniform partition is used for the temporal discretization, and the notations k,  $t_n$  still denotes the time step and time nodes in the partition. The trapezoidal rule is applied to approximate the integral in (4.9).

Thus the numerical method for Problem 8 is defined as follows.

**Problem 9.** Find a discrete displacement  $\boldsymbol{u}^{kh} := \{\boldsymbol{u}_n^{kh}\}_{n=0}^N \subset K^h$  such that for  $0 \le n \le N$ ,

$$\left(\mathcal{A}\boldsymbol{\varepsilon}(\boldsymbol{u}_{n}^{kh}),\boldsymbol{\varepsilon}(\boldsymbol{v}^{h})-\boldsymbol{\varepsilon}(\boldsymbol{u}_{n}^{kh})\right)_{\mathcal{H}}+\left(k\sum_{i=0}^{n}{}^{\prime}\mathcal{B}(t_{n}-t_{i})\boldsymbol{\varepsilon}(\boldsymbol{u}_{i}^{kh}),\boldsymbol{\varepsilon}(\boldsymbol{v}^{h})-\boldsymbol{\varepsilon}(\boldsymbol{u}_{n}^{kh})\right)_{\mathcal{H}} +\int_{\Gamma_{3}}j_{v}^{0}(\boldsymbol{u}_{n,v}^{kh};\boldsymbol{v}_{v}^{h}-\boldsymbol{u}_{n,v}^{kh})d\Gamma \geq \langle \boldsymbol{f}_{n},\boldsymbol{v}^{h}-\boldsymbol{u}_{n}^{kh}\rangle_{V^{*}\times V} \forall \boldsymbol{v}^{h}\in K^{h}.$$
(4.12)

By the same kind of derivation used in proving Theorem 5, we have

$$\max_{0 \le n \le N} \|\boldsymbol{u}_{n} - \boldsymbol{u}_{n}^{kh}\|_{V} \le C \max_{0 \le n \le N} \left\{ \inf_{\boldsymbol{v}^{h} \in K^{h}} \left( \|\boldsymbol{u}_{n} - \boldsymbol{v}^{h}\|_{V} + \|\boldsymbol{u}_{n,\nu} - \boldsymbol{v}_{\nu}^{h}\|_{L^{2}(\Gamma_{3})}^{\frac{1}{2}} + |E_{n}(\boldsymbol{v} - \boldsymbol{u}_{n})|^{\frac{1}{2}} \right) + \inf_{\boldsymbol{v} \in K} |E_{n}(\boldsymbol{v} - \boldsymbol{u}_{n}^{kh})|^{\frac{1}{2}} \right\}$$

$$+ Ck^{2} \|\boldsymbol{u}\|_{W^{2,\infty}(I;V)}, \qquad (4.13)$$

where

$$E_n(\boldsymbol{v}) = \langle A\boldsymbol{u}_n, \boldsymbol{v} \rangle + \langle S_n \boldsymbol{u}, \boldsymbol{v} \rangle + \int_{\Gamma_3} j_{\boldsymbol{v}}^0(\boldsymbol{u}_{n,\boldsymbol{v}}; \boldsymbol{v}_{\boldsymbol{v}}) d\Gamma - \langle \boldsymbol{f}_n, \boldsymbol{v} \rangle.$$
(4.14)

For solution regularity, we assume  $\boldsymbol{u} \in W^{2,\infty}_{loc}(\mathbb{R}_+; V) \cap C(\mathbb{R}_+; H^2(\Omega; \mathbb{R}^d)), \sigma \boldsymbol{v} \in C(\mathbb{R}_+; L^2(\Gamma_3; \mathbb{R}^d)), \boldsymbol{u}|_{\Gamma_{3,l}} \in C(\mathbb{R}_+; H^2(\Gamma_3; \mathbb{R}^d))$  $(\Gamma_{3,l}; \mathbb{R}^d)$ ) and  $g|_{\Gamma_{3,l}} \in H^2(\Gamma_{3,l}), 1 \le l \le l_3$ . Note that if  $\mathcal{A}$  and  $\mathcal{B}$  are suitably smooth, then we can derive  $\sigma \in C(\mathbb{R}_+; H^1(\Omega; \mathbb{S}^d))$  from  $C(\mathbb{R}_+; H^2(\Omega; \mathbb{R}^d))$ . In addition,

the condition  $\sigma \in C(\mathbb{R}_+; H^1(\Omega; \mathbb{S}^d))$  implies  $\sigma \nu \in C(\mathbb{R}_+; L^2(\Gamma_3; \mathbb{R}^d))$ .

Now we bound the residual term  $E_n(\mathbf{v})$ . Define a subset of K,

 $\widetilde{K} = \{ \boldsymbol{\nu} \in C^{\infty}(\overline{\Omega})^d | \boldsymbol{\nu} = \boldsymbol{0} \text{ on } \Gamma_1 \cup \Gamma_3 \}.$ 

Pointwise equations will be shown from the weak formulation (4.7) and the arguments in [21, Section 8.1]. Define

$$\boldsymbol{\sigma}(t) = \mathcal{A}\boldsymbol{\varepsilon}(\boldsymbol{u}(t)) + \int_0^t \mathcal{B}(t-s)\boldsymbol{\varepsilon}(\boldsymbol{u}(s))ds \text{ in } \boldsymbol{\Omega}.$$

Let  $\tilde{v}$  be an arbitrary function from the subset  $\tilde{K}$ , and take  $v = u \pm \tilde{v}$  in (4.7) to get

$$\left(\boldsymbol{\sigma}(t), \boldsymbol{\varepsilon}(\widetilde{\boldsymbol{v}})\right)_{\mathcal{H}} = \langle \boldsymbol{f}(t), \widetilde{\boldsymbol{v}} \rangle_{V^* \times V} \ \forall \ \widetilde{\boldsymbol{v}} \in \widetilde{K}.$$

The following two equations are derived from the above relation:

$$\operatorname{Div} \boldsymbol{\sigma}(t) + \boldsymbol{f}_0(t) = \boldsymbol{0} \text{ a.e. in } \Omega, \tag{4.15}$$

$$\boldsymbol{\sigma}(t)\boldsymbol{\nu} = \boldsymbol{f}_2(t) \text{ a.e. on } \boldsymbol{\Gamma}_2. \tag{4.16}$$

Taking  $t = t_n$  in (4.15) and (4.16), we get

$$\operatorname{Div} \boldsymbol{\sigma}(t_n) + \boldsymbol{f}_0(t_n) = \boldsymbol{0} \text{ a.e. in } \Omega, \tag{4.17}$$

$$\boldsymbol{\sigma}(t_n)\boldsymbol{\nu} = \boldsymbol{f}_2(t_n) \text{ a.e. on } \boldsymbol{\Gamma}_2. \tag{4.18}$$

Multiplying (4.17) by  $\boldsymbol{v}$  and integrating over  $\Omega$ , we obtain

$$\int_{\Omega} \operatorname{Div} \boldsymbol{\sigma}(t_n) \cdot \boldsymbol{\nu} dx + \int_{\Omega} \boldsymbol{f}_0(t_n) \cdot \boldsymbol{\nu} dx = 0.$$
(4.19)

Performing an integration by parts, we get

$$\int_{\Omega} \boldsymbol{\sigma}(t_n) \cdot \boldsymbol{\varepsilon}(\boldsymbol{v}) d\boldsymbol{x} = \int_{\Omega} \boldsymbol{f}_0(t_n) \cdot \boldsymbol{v} d\boldsymbol{x} + \int_{\Gamma} \boldsymbol{\sigma}(t_n) \boldsymbol{v} \cdot \boldsymbol{v} d\Gamma.$$
(4.20)

Therefore

$$\int_{\Omega} \mathcal{A}\boldsymbol{\varepsilon}(\boldsymbol{u}_{n}) \cdot \boldsymbol{\varepsilon}(\boldsymbol{v}) dx + \int_{\Omega} \int_{0}^{t_{n}} \mathcal{B}(t_{n} - s)\boldsymbol{\varepsilon}(\boldsymbol{u}(s)) ds \cdot \boldsymbol{\varepsilon}(\boldsymbol{v}) dx$$
$$= \int_{\Omega} \boldsymbol{f}_{0}(t_{n}) \cdot \boldsymbol{v} dx + \int_{\Gamma_{2}} \boldsymbol{f}_{2}(t_{n}) \cdot \boldsymbol{v} d\Gamma + \int_{\Gamma_{3}} \boldsymbol{\sigma}(t_{n}) \boldsymbol{v} \cdot \boldsymbol{v} d\Gamma,$$

i.e.,

$$\langle A\boldsymbol{u}_n, \boldsymbol{v} \rangle + \langle S_n \boldsymbol{u}, \boldsymbol{v} \rangle - \langle \boldsymbol{f}_n, \boldsymbol{v} \rangle = \int_{\Gamma_3} \boldsymbol{\sigma}(t_n) \boldsymbol{v} \cdot \boldsymbol{v} d\Gamma.$$
(4.21)

Therefore

$$E_n(\mathbf{v}) = \int_{\Gamma_3} \boldsymbol{\sigma}(t_n) \mathbf{v} \cdot \mathbf{v} d\Gamma + \int_{\Gamma_3} j_{\nu}^0(u_{n,\nu}; v_{\nu}) d\Gamma.$$
(4.22)

Thus

$$|E_n(\mathbf{v})| \leq |\int_{\Gamma_3} \boldsymbol{\sigma}(t_n) \boldsymbol{v} \cdot \boldsymbol{v} d\Gamma| + |\int_{\Gamma_3} j_{\boldsymbol{\nu}}^0(u_{n,\boldsymbol{\nu}}; v_{\boldsymbol{\nu}}) d\Gamma|.$$

Combining with  $\sigma v \in C(\mathbb{R}_+; L^2(\Gamma_3; \mathbb{R}^d))$ , we have

$$|E_n(\mathbf{v})| \le C \|\mathbf{v}\|_{L^2(\Gamma_3;\mathbb{R}^d)}.$$
(4.23)

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Fig. 1. Initial configuration of the contact problem.

Then we derive the bounds

$$|E_n(\mathbf{v}^h - \mathbf{u}_n)| \le C \|\mathbf{u}_n - \mathbf{v}^h\|_{L^2(\Gamma_3; \mathbb{R}^d)},\tag{4.24}$$

$$|E_n(\mathbf{v} - \mathbf{u}_n^{\kappa n})| \le C \|\mathbf{v} - \mathbf{u}_n^{\kappa n}\|_{L^2(\Gamma_3; \mathbb{R}^d)}.$$
(4.25)

Then from (4.13).

$$\max_{0 \le n \le N} \|\boldsymbol{u}_{n} - \boldsymbol{u}_{n}^{kh}\|_{V} \le C \max_{0 \le n \le N} \left\{ \inf_{\boldsymbol{v}^{h} \in K^{h}} \left( \|\boldsymbol{u}_{n} - \boldsymbol{v}^{h}\|_{V} + \|\boldsymbol{u}_{n} - \boldsymbol{v}^{h}\|_{L^{2}(\Gamma_{3};\mathbb{R}^{d})}^{\frac{1}{2}} \right) + \inf_{\boldsymbol{v} \in K} \|\boldsymbol{v} - \boldsymbol{u}_{n}^{kh}\|_{L^{2}(\Gamma_{3};\mathbb{R}^{d})}^{\frac{1}{2}} \right\} + Ck^{2} \|\boldsymbol{u}\|_{W^{2,\infty}(l;V)}.$$

$$(4.26)$$

Now we focus on the estimate of the third term on the right hand of (4.26), which can be solved using the arguments presented in [8, Section 4].

Suppose the contact boundary  $\Gamma_3$  is smooth. For  $\mathbf{x} \in \Gamma_3$ , the unit outward normal  $\mathbf{v}(\mathbf{x})$  is extended to  $\overline{\mathbf{v}}$ , where  $\overline{\mathbf{v}} \in H^1(\Omega)$ and  $\overline{v}|_{\Gamma_3} = v(\mathbf{x})$ . Let  $\phi(\mathbf{x})$  be a smooth cut-off function defined by  $\phi(\mathbf{x}) = \mathbf{0}$  outside the small neighborhood of  $\Gamma_3$  and  $\phi(\mathbf{x}) = \mathbf{1}$ near  $\Gamma_3$ . Let  $\overline{v}$  multiplying  $\phi(\mathbf{x})$ , we obtain a function  $\widetilde{v} \in H^1(\Omega; \mathbb{R}^d)$  with  $\widetilde{v}|_{\Gamma_3} = v$ . Assume  $\widetilde{g} \in H^1(\Omega; \mathbb{R}^d)$ ,  $\widetilde{g}|_{\Gamma_3} = g$  and g is continuous. Define

$$\boldsymbol{v} = \boldsymbol{u}_n^{kh} + (\min\{\widetilde{g}, \boldsymbol{u}_n^{kh} \cdot \widetilde{\boldsymbol{v}}\} - \boldsymbol{u}_n^{kh} \cdot \widetilde{\boldsymbol{v}})\widetilde{\boldsymbol{v}}$$

Thus on  $\Gamma_3$ , we have  $\boldsymbol{v}_{\tau} = \boldsymbol{u}_{n,\tau}^{kh}$ ,  $v_{\nu} = \min\{g, u_{n,\nu}^{kh}\}$ . Therefor we get  $v_{\nu} \leq g$  on  $\Gamma_3$ , i.e.,  $\boldsymbol{v} \in K$ . We observe that

$$\|\boldsymbol{v} - \boldsymbol{u}_n^{kh}\|_{L^2(\Gamma_3;\mathbb{R}^d)} = \|v_v - u_{n,v}^{kh}\|_{L^2(\Gamma_3)}.$$
(4.27)

Note that  $v_v = \min\{g, u_{n,v}^{kh}\}$ , thus on  $\Gamma_3, 0 \le u_{n,v}^{kh}(\mathbf{x}) - v_v(\mathbf{x})$ . Since  $u_{n,v}^{kh}(\mathbf{x}) \le g(\mathbf{x})$  at nodes  $\mathbf{x}$  on  $\overline{\Gamma_3}$ , we obtain  $u_{n,v}^{kh}(\mathbf{x}) \le g(\mathbf{x})$ .

 $\Pi^h g(\mathbf{x}), \ \mathbf{x} \in \Gamma_3$ , where  $\Pi^h g(\mathbf{x})$  is the linear finite element interpolate of g on  $\overline{\Gamma_3}$ . Then we consider two cases on  $\Gamma_3$ : if  $u_{n,\nu}^{kh} \leq g$ , then  $v_{\nu} = u_{n,\nu}^{kh}$ ; if  $u_{n,\nu}^{kh} > g$ , then  $v_{\nu} = g$ , and thus  $0 \leq u_{n,\nu}^{kh}(\mathbf{x}) - v_{\nu}(\mathbf{x}) \leq g$ .  $\Pi^{h}g(\mathbf{x}) - g(\mathbf{x})$ . Therefore

$$\|v_{\nu} - u_{n,\nu}^{kh}\|_{L^{2}(\Gamma_{3})} \le \|\Pi^{h}g - g\|_{L^{2}(\Gamma_{3})}.$$
(4.28)

Combining (4.27) and (4.28), we have

$$\|\boldsymbol{v} - \boldsymbol{u}_n^n\|_{L^2(\Gamma_3;\mathbb{R}^d)} \le \|\Pi^n g - g\|_{L^2(\Gamma_3)}.$$
(4.29)

Then applying finite element interpolation theory [23,24], we have the optimal order error bound for Problem 9

$$\max_{0 \le n \le N} \| \boldsymbol{u}_n - \boldsymbol{u}_n^{\kappa n} \|_V \le C(h + k^2).$$
(4.30)

Thus in spatial mesh size the method is of first-order, and in the time step it is of second-order.

#### 5. A numerical example

In this section we report a numerical example for Problem 7. The physical setting of the problem is depicted in Fig. 1. Let  $\Omega = (0, L_1) \times (0, L_2)$  be a rectangle with a boundary  $\Gamma$  which is divided into three parts

$$\Gamma_1 = \{0\} \times (0, L_2), \ \Gamma_2 = \{L_1\} \times (0, L_2) \cup [0, L_1] \times \{L_2\}, \ \Gamma_3 = [0, L_1] \times \{0\}.$$



**Fig. 2.** Normal displacements on  $\Gamma_3$  at different time steps.

The elasticity tensor A satisfies

$$(\mathcal{A}\tau)_{ij} = \frac{E\kappa}{1-\kappa^2} (\tau_{11}+\tau_{22})\delta_{ij} + \frac{E}{1+\kappa} \tau_{ij}, \ 1 \le i,j \le 2.$$
(5.1)

The notation *E* denotes the Young modulus,  $\kappa$  denotes the Poisson ratio of the material, and  $\delta_{ij}$  denotes the Kronecker symbol. The relaxation tensor  $\mathcal{B}(s) = (0.5 + s)^3 \mathbb{I}$ , where  $\mathbb{I}$  is identity matrix. For a given value  $S \ge 0$ , the function  $j_{\nu}(\cdot)$  is defined by

$$j_{\nu}(u_{\nu}) = S \int_{0}^{|u_{\nu}|} \mu(s) ds,$$
(5.2)

and the function  $\mu(\cdot)$  defined by

$$\mu(s) = (a-b)e^{-\alpha s} + b, \tag{5.3}$$

with  $a \ge b > 0$  and  $\alpha > 0$ . The following parameters are used in numerical experiments:

$$L_1 = 2 \text{ m}, \quad L_2 = 1 \text{ m}, \quad E = 1 \text{ N/m}^2, \quad \kappa = 0.3,$$
  
 $\alpha = 100, \quad a = 0.04, \quad b = 0.02, \quad S = 0.1 \text{ N}, \quad g = 0.2 \text{ m},$ 

and

$$\mathbf{f}_2 = \begin{cases} (0,0) \ N/m & \text{on } \{L_1\} \times (0,L_2) \\ (0,-0.1(1-e^{-2t})\mathbf{x}) \ N/m & \text{on } [0,L_1] \times \{L_2\}. \end{cases}$$
(5.4)

To numerically solve the sample problem, uniform rectangular finite element partitions are introduced, with the intervals  $[0, L_1]$  and  $[0, L_2]$  being divided into 1/h equal parts, and the corresponding continuous bilinear rectangular finite element space for  $V^h$  is used.

The normal displacements of the deformable body on  $\Gamma_3$  at different time steps are showed in Fig. 2. To compute the numerical solution errors, the numerical solution with  $h = k = \frac{1}{256}$  is used as the "reference" solution. The numerical convergence orders in  $H^1$  norm are reported in Tables 1 and 2 which show the first order convergence for spatial discretization and second order for time discretization, and confirm the previous theoretical results.

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#### Table 1

Convergence orders of the errors with fixed time step size.

h	k	$\left\ \boldsymbol{u}(\cdot,T)-\boldsymbol{u}_{N}^{kh}\right\ _{1}$	Order
1/8	1/256	4.049067247799612e-03	_
1/16	1/256	9.899921240195822e-04	1.31
1/32	1/256	4.290570202713140e-04	1.19
1/64	1/256	2.416076212233180e-04	1.13

#### Table 2

Convergence orders of the errors with fixed spatial step size.

h	k	$\left\ \boldsymbol{u}(\cdot,T)-\boldsymbol{u}_{N}^{kh}\right\ _{1}$	Order
1/256	1/4	7.044095264494761e-03	_
1/256	1/8	2.831608699747937e-03	2.03
1/256	1/12	1.244870264157188e-03	2.06
1/256	1/16	5.692251463968669e-04	2.00

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