

# A NEW APPROACH TO THE NONSINGULAR CUBIC BINARY MOMENT PROBLEM

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ABSTRACT. We present an alternative solution to nonsingular cubic moment problems, using techniques that are expected to be useful for higher-degree truncated moment problems. In particular, we apply the theory of recursively determinate moment matrices to deal with a case of rank-increasing moment matrix extensions.

## 1. INTRODUCTION

Given a doubly indexed finite sequence of real numbers  $\beta \equiv \beta^{(m)} = \{\beta_{00}, \beta_{10}, \beta_{01}, \dots, \beta_{m,0}, \beta_{m-1,1}, \dots, \beta_{1,m-1}, \beta_{0,m}\}$  with  $\beta_{00} > 0$ , the *truncated real moment problem* (TRMP) entails seeking necessary and sufficient conditions for the existence of a positive Borel measure  $\mu$  supported in the real plane  $\mathbb{R}^2$  such that

$$\beta_{ij} = \int x^i y^j d\mu \quad (i, j \in \mathbb{Z}_+, 0 \leq i + j \leq m).$$

When such a measure exists, we say that  $\mu$  is a *representing measure* for  $\beta$  and that the TRMP is *soluble*.

There is a parallel *truncated complex moment problem* (TCMP), for a finite sequence of complex numbers  $\gamma \equiv \gamma^{(m)} : \gamma_{00}, \gamma_{01}, \gamma_{10}, \dots, \gamma_{0,m}, \gamma_{1,m-1}, \dots, \gamma_{m-1,1}, \gamma_{m,0}$ , with  $\gamma_{00} > 0$  and  $\gamma_{ji} = \bar{\gamma}_{ij}$ . Here TCMP consists of finding a positive Borel measure  $\mu$  supported in the complex plane  $\mathbb{C}$  such that  $\gamma_{ij} = \int \bar{z}^i z^j d\mu$  ( $0 \leq i + j \leq m$ ). It is well-known that TRMP and TCMP are equivalent for an even integer  $m$  (cf. [4, Proposition 1.12]), and hence any techniques developed for TCMP are transferable to TRMP. Both problems are simply referred to as the *truncated moment problem* (TMP).

In a series of papers, for the case when  $m = 2d$ , the first named author and L.A. Fialkow found solutions for various truncated moment problems; for instance, they obtained complete solutions for  $m = 2$  and  $m = 4$  (cf. [1, 4, 10, 9]). Some solutions are based on matrix positivity and extension, combined with a “functional calculus” (to be briefly discussed in Section 2) for the columns of the associated moment matrix. This matrix is defined as follows. For a real moment sequence  $\beta^{(2d)}$  of even degree, the *moment matrix*  $\mathcal{M}(d) \equiv \mathcal{M}(d)(\beta^{(2d)})$  is given

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by

$$\mathcal{M}(d)(\beta^{(2d)}) := (\beta_{\mathbf{i}+\mathbf{j}})_{\mathbf{i}, \mathbf{j} \in \mathbb{Z}_+^2: |\mathbf{i}|, |\mathbf{j}| \leq 2d}.$$

If we label the columns of  $\mathcal{M}(d)$  with the degree lexicographical order,  $1, X, Y, X^2, XY, Y^2, \dots, X^d, \dots, Y^d$ , we can then use the functional calculus for columns of  $\mathcal{M}(d)$ , introduced in [1]. The moment matrix  $\mathcal{M}(d)$  is Hankel by rectangular blocks; for instance,

$$\mathcal{M}(2) \equiv \begin{pmatrix} \beta_{00} & | & \beta_{10} & \beta_{01} & | & \beta_{20} & \beta_{11} & \beta_{02} \\ \hline \beta_{10} & | & \beta_{20} & \beta_{11} & | & \beta_{30} & \beta_{21} & \beta_{12} \\ \beta_{01} & | & \beta_{11} & \beta_{02} & | & \beta_{21} & \beta_{12} & \beta_{03} \\ \hline \beta_{20} & | & \beta_{30} & \beta_{21} & | & \beta_{40} & \beta_{31} & \beta_{22} \\ \beta_{11} & | & \beta_{21} & \beta_{12} & | & \beta_{31} & \beta_{22} & \beta_{13} \\ \beta_{02} & | & \beta_{12} & \beta_{03} & | & \beta_{22} & \beta_{13} & \beta_{04} \end{pmatrix}. \quad (1.1)$$

When  $m = 2d + 1$ , a general solution to partial cases of the TMP can be found in [11] and [13] as well as a solution to the *truncated matrix moment problem*; a solution to the *cubic complex moment problem* (when  $m = 3$ ) was given in [11]. We know from [4, Proposition 1.12] that the complex and real truncated moment problems are equivalent, in the sense that there exists a bridging map that allows one to translate the hypotheses and conclusions for TCMP into similar hypotheses and conclusions for TRMP, and vice versa. Our approach here, however, focuses on the real cubic moment problem, and analyzes it in its own right.

In this note we consider cubic *real* moment problems and present an alternative solution to the “nonsingular” case (i.e.,  $\mathcal{M}(1)$  invertible; see Section 3 for the formal definition). Our idea is to extend the initial data  $\beta^{(3)}$  to an even-degree  $\beta^{(4)}$ , for which the associated moment matrix  $\mathcal{M}(2)$  has rank 3 or 4. We then prove that  $\mathcal{M}(2)$  is a flat extension of  $\mathcal{M}(1)$ , or is a flat extension of a  $4 \times 4$  submatrix, or admits a flat extension  $\mathcal{M}(3)$ . In all three cases we find a finitely atomic representing measure for  $\beta^{(3)}$ .

We anticipate that the present work will contribute to our understanding of higher-degree moment problems, beginning with the quintic moment problem. We also expect that solutions to odd-degree moment problems will be applied to solve the subnormal completion problem studied in [8].

## 2. PRELIMINARIES

When we build a moment matrix  $\mathcal{M}(2)$  out of a cubic finite sequence, the lower right-hand  $3 \times 3$  block will include all quartic moments, which will need to remain undefined. To obtain our main results, we will choose appropriate quartic moments and show that  $\mathcal{M}(2)$  has a representing measure. In order to describe this process in detail, we need to review basic TMP notation and results pertaining to the even-degree case.

**Necessary Conditions.** In order to discuss basic necessary conditions for the existence of a measure, let  $\mu$  be a representing measure of the even-degree moment sequence  $\beta \equiv \beta^{(2d)}$ . First, we recall that

$$0 \leq \int |p(x, y)|^2 d\mu = \sum_{i,j,k,l} a_{ij} a_{kl} \int x^{i+l} y^{j+k} d\mu = \sum_{i,j,k,l} a_{ij} a_{kl} \beta_{i+l} \beta_{j+k}$$

if and only if  $\mathcal{M}(d) \geq 0$ .

Let  $\mathcal{P}_k$  denote the set of bivariate polynomials in  $\mathbb{R}[x, y]$  whose degree is at most  $k$  and let  $\mathcal{C}_{\mathcal{M}(d)}$  denote the column space of  $\mathcal{M}(d)$ . For  $k \leq d$ , we now define an assignment from  $\mathcal{P}_k$  to  $\mathcal{C}_{\mathcal{M}(d)}$ ; given a polynomial  $p(x, y) \equiv \sum_{ij} a_{ij} x^i y^j$ , we let  $p(X, Y) := \sum_{ij} a_{ij} X^i Y^j$  (so that  $p(X, Y) \in \mathcal{C}_{\mathcal{M}(d)}$ ), which defines the above mentioned functional calculus. We also let  $\mathcal{Z}(p)$  denote the zero set of  $p$  and define the *algebraic variety* of  $\beta$  by

$$\mathcal{V} \equiv \mathcal{V}(\beta) := \bigcap_{p(x,y)=0, \deg p \leq d} \mathcal{Z}(p). \quad (2.1)$$

If  $\hat{p}$  denotes the column vector of coefficients of  $p$ , then we know that  $p(X, Y) = \mathcal{M}(d)\hat{p}$ ; as a consequence,  $p(X, Y) = \mathbf{0}$  if and only if  $\hat{p} \in \ker \mathcal{M}(d)$ . Another necessary condition we will use is  $\text{supp } \mu \subseteq \mathcal{V}(\beta)$  and  $r := \text{rank } \mathcal{M}(d) \leq \text{card } \text{supp } \mu \leq v := \text{card } \mathcal{V}$ ; this condition is called the *variety condition* [1]. In addition, if  $p$  is any polynomial of degree at most  $2d$  such that  $p|_{\mathcal{V}} \equiv 0$ , then the *Riesz functional*  $\Lambda$  must satisfy  $\Lambda(p) := \int p d\mu = 0$ , which is referred to as *Consistency* of the moment sequence. The main results in [6] state that the above mentioned conditions together with Consistency are sufficient for solubility in the extremal case ( $r = v$ ); moreover, in [6] the authors showed that Consistency cannot be replaced by the weaker condition that  $\mathcal{M}(d)$  is *recursively generated*, (RG), that is,  $p(X, Y) = \mathbf{0} \Rightarrow (pq)(X, Y) = \mathbf{0}$ , for each polynomial  $q$  with  $\deg(pq) \leq d$ .

In summary, positive semidefiniteness alone is sufficient to solve the quadratic moment problem ( $d = 1$ ) [1]. However, for  $d > 1$ , the solubility of TMP requires more. For instance, the solution of the quartic moment problem ( $d = 2$ ) requires positive semidefiniteness, the variety condition, and the (RG) property (this last property requires that the moment matrix be recursively generated; cf. [1], [4], [10]).

**Flat Extensions.** We recall that  $\mathcal{M}(d)$  is said to be *flat* if  $\text{rank } \mathcal{M}(d) = \text{rank } \mathcal{M}(d-1)$ , that is,  $\mathcal{M}(d)$  is a rank-preserving positive extension of  $\mathcal{M}(d-1)$ ; in this case,  $\mathcal{M}(d)$  has a unique rank  $\mathcal{M}(d)$ -atomic measure. Furthermore, it is known that if  $\mathcal{M}(d)$  has a positive extension  $\mathcal{M}(d+k)$  for some  $k \in \mathbb{Z}_+$ , which in turn admits a flat extension  $\mathcal{M}(d+k+1)$ , then  $\beta$  has a rank  $\mathcal{M}(d+k)$ -atomic measure [3, Theorem 1.5]. This result is referred to as the Flat Extension Theorem; it is probably the most efficient, concrete solution to TMP, even though the construction of an extension is usually difficult for a high-degree TMP.

We will use the Flat Extension Theorem in the proof of our main result; thus, we need to briefly describe the process for building a flat extension. Since a moment matrix extension  $\mathcal{M}(d+1)$  of  $\mathcal{M}(d)$  can be written as  $\mathcal{M}(d+1) =$

$\begin{pmatrix} \mathcal{M}(d) & B(d+1) \\ B(d+1)^* & C(d+1) \end{pmatrix}$ , for some rectangular matrices  $B(d+1)$  and  $C(d+1)$ , we can adapt a classical result given by J.L. Smul'jan in the search for a positive  $\mathcal{M}(d+1)$ .

**Theorem 2.1.** (*Smul'jan's Theorem [14]*) *Let  $A, B, C$  be matrices of complex numbers, with  $A$  and  $C$  square matrices. Then*

$$\tilde{A} := \begin{pmatrix} A & B \\ B^* & C \end{pmatrix} \geq 0 \iff \begin{cases} A \geq 0 \\ B = AW \text{ (for some } W) \\ C \geq W^*AW \end{cases} .$$

Moreover,  $\text{rank } \tilde{A} = \text{rank } A \iff C = W^*AW$ .

*Remark 2.2.* When the extension  $\tilde{A}$  in Theorem 2.1 has the same rank as  $A$ , we say that  $\tilde{A}$  is a flat extension of  $A$ . Besides satisfying this theorem, an extension  $\mathcal{M}(d+1)$  must maintain the moment matrix structure, that is, the  $C$ -block must be Hankel. This condition makes generating flat extensions quite difficult in many instances.

We now discuss how we can find an explicit formula for a representing measure. Suppose  $\mathcal{M}(d)$  admits a positive extension  $\mathcal{M}(d+k)$  for some  $k \in \mathbb{Z}_+$  that has a flat extension  $\mathcal{M}(d+k+1)$ . Thus,  $\beta$  has a rank  $\mathcal{M}(d+k)$ -atomic measure  $\mu$  and let  $r := \text{rank } \mathcal{M}(d+k)$ . The flat extension theorem says that the algebraic variety  $\mathcal{V}$  of  $\mathcal{M}(d+k+1)$  consists of exactly  $r$  points and we may write  $\mathcal{V} = \{(x_1, y_1), \dots, (x_r, y_r)\}$ . Denote the Vandermonde matrix  $V$  as

$$V = \begin{pmatrix} 1 & x_1 & y_1 & x_1^2 & x_1 y_1 & y_1^2 & \cdots & x_1^{d+k} & \cdots & y_1^{d+k} \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 1 & x_r & y_r & x_r^2 & x_r y_r & y_r^2 & \cdots & x_r^{d+k} & \cdots & y_r^{d+k} \end{pmatrix} \quad (2.2)$$

If  $\mathcal{B} := \{\mathbf{t}_1, \dots, \mathbf{t}_r\}$  is the basis for the column space of  $\mathcal{M}(d+k)$  and if  $V_{\mathcal{B}}$  is the submatrix of  $V$  with columns labeled as in  $\mathcal{B}$ , then we can find the densities by solving:

$$V_{\mathcal{B}}^T (\rho_1 \ \rho_2 \ \cdots \ \rho_r)^T = (\Lambda(\mathbf{t}_1) \ \Lambda(\mathbf{t}_2) \ \cdots \ \Lambda(\mathbf{t}_r))^T. \quad (2.3)$$

Finally, we have  $\mu = \sum_{k=1}^r \rho_k \delta_{(x_k, y_k)}$ .

**Degree-One Transformations.** We briefly review a tool that will allow us to convert the given moment problem into a simpler one; this tool is known as the invariance of moment problems under degree-one transformations. The complex version is provided in [4]; we adapt the notation in [4] to obtain a real version.

For  $a, b, c, d, e, f \in \mathbb{R}$  with  $bf \neq ce$ , let  $\Psi(x, y) \equiv (\Psi_1(x, y), \Psi_2(x, y)) := (a + bx + cy, d + ex + fy)$  for  $x, y \in \mathbb{R}$ . If  $\Lambda_{\beta}$  denotes the Riesz functional associated with  $\beta$ , then given  $\beta \equiv \beta^{(2d)}$  we build a new (equivalent) moment sequence  $\tilde{\beta} \equiv \tilde{\beta}^{(2d)} \equiv \{\tilde{\beta}_{ij}\}$  given by  $\tilde{\beta}_{ij} := \Lambda_{\beta}(\Psi_1^i \Psi_2^j)$  ( $0 \leq i + j \leq 2d$ ). We immediately check that  $\Lambda_{\tilde{\beta}}(p) = \Lambda_{\beta}(p \circ \Psi)$  for every  $p \in \mathcal{P}_d$ .

**Proposition 2.3.** [4, cf. Proposition 1.7] (Invariance under degree-one transformations) *Let  $\mathcal{M}(d)$  and  $\tilde{\mathcal{M}}(d)$  be the moment matrices associated with  $\beta$  and  $\tilde{\beta}$ , resp., and let  $J\hat{p} := \widehat{p \circ \Psi}$  ( $p \in \mathcal{P}_d$ ). The following statements hold:*

- (i)  $\tilde{\mathcal{M}}(d) = J^T \mathcal{M}(d) J$ ;
- (ii)  $J$  is invertible;
- (iii)  $\tilde{\mathcal{M}}(d) \geq 0 \Leftrightarrow \mathcal{M}(d) \geq 0$ ;
- (iv)  $\text{rank } \tilde{\mathcal{M}}(d) = \text{rank } \mathcal{M}(d)$ ;
- (v) The formula  $\mu = \tilde{\mu} \circ \Psi$  establishes a one-to-one correspondence between the sets of representing measures for  $\beta$  and  $\tilde{\beta}$ , which preserves measure class and cardinality of the support; moreover,  $\varphi(\text{supp } \mu) = \text{supp } \tilde{\mu}$ ;
- (vi)  $\mathcal{M}(d)$  admits a flat extension if and only if  $\tilde{\mathcal{M}}(d)$  admits a flat extension.

We will now apply Proposition 2.3 to a cubic real binary moment sequence  $\beta \equiv \beta^{(3)} : \{\beta_{00}, \beta_{10}, \beta_{01}, \beta_{20}, \beta_{11}, \beta_{02}, \beta_{30}, \beta_{21}, \beta_{12}, \beta_{03}\}$  with  $\beta_{00} > 0$ . Our strategy is to enlarge  $\beta^{(3)}$  to  $\beta^{(4)}$  by adding new undetermined moments of degree 4; this extended finite sequence has an associated moment matrix  $\mathcal{M}(2)$ . As we will see in the sequel, it is enough to consider the case when  $\mathcal{M}(2)$  is “normalized,” that is,  $\mathcal{M}(1)$  is the identity matrix. For, the case when  $\mathcal{M}(1)$  is singular can be dealt with easily using the results in [1] and [2]. We thus assume that  $\beta_{00} = 1$  and that the principal  $2 \times 2$  and  $3 \times 3$  minors of  $\mathcal{M}(1)$ ,  $d_2$  and  $d_3$ , resp., are strictly positive. A calculation using *Mathematica* [15] reveals that

$$\begin{aligned} d_2 &= -\beta_{10}^2 + \beta_{20} \\ d_3 &= -\beta_{02}\beta_{10}^2 + 2\beta_{01}\beta_{10}\beta_{11} - \beta_{11}^2 - \beta_{01}^2\beta_{20} + \beta_{02}\beta_{20}. \end{aligned}$$

Consider now the degree-one transformation

$$\Psi(x, y) \equiv (a + bx + cy, d + ex + fy),$$

where  $a := \frac{\beta_{01}\beta_{20} - \beta_{10}\beta_{11}}{\sqrt{d_2d_3}}$ ,  $b := \frac{\beta_{11} - \beta_{01}\beta_{10}}{\sqrt{d_2d_3}}$ ,  $c := -\sqrt{\frac{d_2}{d_3}}$ ,  $d := -\frac{\beta_{10}}{\sqrt{d_2}}$ ,  $e := \frac{1}{\sqrt{d_2}}$ , and  $f := 0$ . Observe that

$$bf - ce = -\sqrt{\frac{1}{d_3}} \neq 0.$$

Through this transformation, and using [4, Proposition 1.7], any positive semi-definite  $\mathcal{M}(2)$  with a nonsingular  $\mathcal{M}(1)$  can be translated to

$$\begin{pmatrix} 1 & 0 & 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & \tilde{\beta}_{30} & \tilde{\beta}_{21} & \tilde{\beta}_{12} \\ 0 & 0 & 1 & \tilde{\beta}_{21} & \tilde{\beta}_{12} & \tilde{\beta}_{03} \\ 1 & \tilde{\beta}_{30} & \tilde{\beta}_{21} & \tilde{\beta}_{40} & \tilde{\beta}_{31} & \tilde{\beta}_{22} \\ 0 & \tilde{\beta}_{21} & \tilde{\beta}_{12} & \tilde{\beta}_{31} & \tilde{\beta}_{22} & \tilde{\beta}_{13} \\ 1 & \tilde{\beta}_{12} & \tilde{\beta}_{03} & \tilde{\beta}_{22} & \tilde{\beta}_{13} & \tilde{\beta}_{04} \end{pmatrix} =: \mathcal{M}[a_0, a_1, a_2, a_3] \quad (2.4)$$

where  $a_i := \tilde{\beta}_{3-i,i}$ .

**Recursively Determinate Moment Problems.** Our approach to the nonsingular cubic moment problem will require a key result from the theory of recursively determinate moment matrices [7], which we now briefly describe. We first recall that a moment matrix  $\mathcal{M}(d)$  is *recursively determinate* if there are column dependence relations in  $\mathcal{M}(d)$  of the form

$$X^n = p(X, Y) \quad (p \in \mathcal{P}_{n-1}); \quad (2.5)$$

$$Y^m = q(X, Y) \quad (q \in \mathcal{P}_m, m \leq n, \text{ and } q \text{ has no } y^m \text{ term}), \quad (2.6)$$

where  $\mathcal{P}_k$  denotes the subspace of polynomials in  $\mathbb{R}[x, y]$  whose degree is less than or equal to  $k$ . One of the main results in [7] follows.

**Lemma 2.4.** ([7, Corollary 2.4], with  $d = n = m = 2$ , so that  $d = n + m - 2$ .) Assume that  $\mathcal{M}(2)$  is positive semidefinite and admits column relations of the form (2.5) and (2.6), with  $n = m = 2$ . Then  $\mathcal{M}(2)$  admits a flat extension  $\mathcal{M}(3)$ .

We recall that, in general, the solubility of a quartic moment problem requires the variety condition. However, Lemma 2.4 says that the variety condition is superfluous if a positive semidefinite  $\mathcal{M}(2)$  with invertible  $\mathcal{M}(1)$  has only two column relations  $X^2 = p(X, Y)$  and  $Y^2 = q(X, Y)$ , where  $p$  and  $q$  are linear polynomials. In such a case,  $\mathcal{M}(2)$  has a flat extension  $\mathcal{M}(3)$ , and therefore a 4-atomic representing measure. It follows that the pair of equations  $x^2 = p(x, y)$  and  $y^2 = q(x, y)$  has exactly 4 common real roots.

### 3. CUBIC BINARY MOMENT PROBLEMS

As we have indicated before, the nontrivial cases of the cubic binary moment problem arise when the submatrix  $\mathcal{M}(1)$  of  $\beta^{(3)}$  is nonsingular. Moreover, as noted in [12] the positive semidefiniteness of  $\mathcal{M}(1)$  is always a necessary condition for the existence of a representing measure. Thus, in the sequel we focus on cubic binary moment problems with  $\mathcal{M}(1)$  positive definite. When this happens we say that  $\beta^{(3)}$  is a *nonsingular* cubic binary moment sequence.

**Main Results.** Using the degree-one transformation introduced in Section 2, if  $\beta^{(3)}$  is a nonsingular cubic binary moment sequence we may always assume, without loss of generality, that  $\beta^{(3)} : \{1, 0, 0, 1, 0, 1, a_0, a_1, a_2, a_3\}$  and we may write

$$\mathcal{M}(2) := \begin{pmatrix} 1 & 0 & 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & a_0 & a_1 & a_2 \\ 0 & 0 & 1 & a_1 & a_2 & a_3 \\ 1 & a_0 & a_1 & \beta_{40} & \beta_{31} & \beta_{22} \\ 0 & a_1 & a_2 & \beta_{31} & \beta_{22} & \beta_{13} \\ 1 & a_2 & a_3 & \beta_{22} & \beta_{13} & \beta_{04} \end{pmatrix}, \quad (3.1)$$

where  $\beta_{40}, \beta_{31}, \beta_{22}, \beta_{13}$  and  $\beta_{04}$  are undetermined new moments. We will prove that the extended  $\beta^{(4)}$  obtained from  $\beta^{(3)}$  by adding the quartic moments  $\beta_{40}, \beta_{31}, \beta_{22}, \beta_{13}$  and  $\beta_{04}$  admits a representing measure, for appropriate choices of the new moments; as a result,  $\beta^{(3)}$  also admits a representing measure  $\mu$ . The smallest cardinality of  $\text{supp } \mu$  will be 3 in some cases, and 4 in others.

Using Theorem 2.1, we first determine under what conditions the extended matrix  $\mathcal{M}(2)$  will be a flat extension of  $\mathcal{M}(1)$ . First, to ensure the positive semidefiniteness of  $\mathcal{M}(2)$ , and if we let  $W := B(2)$  (the upper right-hand  $3 \times 3$  block of  $\mathcal{M}(2)$ ), we see that  $C(2)$  (the lower right-hand  $3 \times 3$  block of  $\mathcal{M}(2)$ ) must satisfy the inequality  $C(2) \geq W^T \mathcal{M}(1) W$ , with equality characterizing flatness. Now,

$$W^T(\mathcal{M}(1))^{-1}W = W^TW = \begin{pmatrix} 1 + a_0^2 + a_1^2 & a_0a_1 + a_1a_2 & 1 + a_0a_2 + a_1a_3 \\ a_0a_1 + a_1a_2 & a_1^2 + a_2^2 & a_1a_2 + a_2a_3 \\ 1 + a_0a_2 + a_1a_3 & a_1a_2 + a_2a_3 & 1 + a_2^2 + a_3^2 \end{pmatrix}. \quad (3.2)$$

Consequently,  $\mathcal{M}(2)$  is a flat extension of  $\mathcal{M}(1)$  if and only if

$$\beta_{40} = 1 + a_0^2 + a_1^2 \quad (3.3)$$

$$\beta_{31} = a_0a_1 + a_1a_2 \quad (3.4)$$

$$\beta_{13} = a_1a_2 + a_2a_3 \quad (3.5)$$

$$\beta_{04} = 1 + a_2^2 + a_3^2, \text{ and} \quad (3.6)$$

$$k := (1 + a_0a_2 + a_1a_3) - (a_1^2 + a_2^2) = 0. \quad (3.7)$$

Condition (3.7) is equivalent to the commutativity of the matrices defined in [12].

We are now ready to prove our first result.

**Theorem 3.1.** *Let  $\beta^{(3)}$  be a nonsingular cubic binary moment sequence, let  $k$  be as in (3.7) and assume that  $k = 0$ . Then  $\beta^{(3)}$  admits a 3-atomic representing measure.*

*Proof.* From the discussion preceding the statement of Theorem 3.1, the new quartic moments  $\beta_{40}$ ,  $\beta_{31}$ ,  $\beta_{13}$  and  $\beta_{04}$  must be defined using (3.3) – (3.7), resp. As for  $\beta_{22}$ , we must use  $1 + a_0a_2 + a_1a_3$ , which in this case equals  $a_1^2 + a_2^2$ , because  $k = 0$ . With these definitions, we easily conclude that  $\mathcal{M}(2)$  is a flat moment matrix extension of  $\mathcal{M}(1)$ , which gives the desired result.  $\square$

When  $k \neq 0$ , it is not possible to select new quartic moments so that  $\mathcal{M}(2)$  is a flat extension of  $\mathcal{M}(1)$ . Therefore, any positive semidefinite moment matrix extension  $\mathcal{M}(2)$  will satisfy  $\text{rank } \mathcal{M}(2) \geq 4$ . Nevertheless, the following theorem shows that it is always possible to choose a set of quartic moments such that  $\text{rank } \mathcal{M}(2) = 4$ . Once those moments have been appropriately chosen, the extended moment matrix  $\mathcal{M}(2)$  will admit a flat extension  $\mathcal{M}(3)$ , and therefore a 4-atomic representing measure for  $\beta^{(4)}$ , which is also a representing measure for the initial data sequence  $\beta^{(3)}$ .

**Theorem 3.2.** *Let  $\beta^{(3)}$  be a nonsingular cubic binary moment sequence, let  $k$  be as in (3.7) and assume that  $k \neq 0$ . Then  $\beta^{(3)}$  admits a 4-atomic representing measure.*

*Proof.* We will divide the proof into two cases:  $k > 0$  and  $k < 0$ .

**Case 1** ( $k > 0$ ). As in the Proof of Theorem 3.1, let  $\beta_{40}$ ,  $\beta_{31}$ ,  $\beta_{13}$  and  $\beta_{04}$  be given by (3.3) – (3.7), resp. Since  $k > 0$ , the positivity of  $\mathcal{M}(2)$  will be

preserved if we let  $\beta_{22} := 1 + a_1a_3 + a_2a_4$ . With this choice of  $\beta_{22}$ , the proposed extended matrix  $\mathcal{M}(2)$  will be a positive semidefinite moment matrix, and such that the block  $C(2)$  differs from  $W^T(\mathcal{M}(1))^{-1}W$  in just the  $(2, 2)$ -entry. As a result,  $\text{rank } \mathcal{M}(2) = 4$ . A simple calculation now reveals that

$$\begin{aligned} X^2 &= 1 + a_0X + a_1Y \\ Y^2 &= 1 + a_2X + a_3Y. \end{aligned}$$

We now know that  $\mathcal{M}(2)$  is positive semidefinite and recursively determinate, and by Lemma 2.4,  $\mathcal{M}(2)$  admits a 4-atomic representing measure; it follows that  $\beta^{(3)}$  also admits a 4-atomic representing measure.

**Case 2** ( $k < 0$ ). Here our strategy is to allow the rank to increase as we transition from  $\mathcal{M}(1)$  to the compression of  $\mathcal{M}(2)$  to the first 4 rows and columns. This requires making the column  $X^2$  *linearly independent* of the columns  $1$ ,  $X$ , and  $Y$  in  $\mathcal{M}(2)$ . It is straightforward to observe that this can be easily accomplished by letting

$$\beta_{40} := 2 + a_1^2 + a_2^2.$$

With this definition in hand, we now postulate that  $\mathcal{M}(2)$  be a flat extension of its compression to the first 4 rows and columns. A calculation using *Mathematica* reveals that one can accomplish this by defining three of the remaining quartic moments as follows:

$$\begin{aligned} \beta_{31} &:= a_1a_2 + a_2a_3, \\ \beta_{22} &:= a_2^2 + a_3^2, \\ \beta_{13} &:= a_2a_3 + a_3a_4. \end{aligned}$$

Having chosen these moments, we now use Theorem 2.1 to determine the remaining quartic moment,  $\beta_{04}$ . Since we wish to make  $\mathcal{M}(2)$  a flat extension of its above mentioned compression, a calculation using *Mathematica* immediately yields

$$\begin{aligned} \beta_{04} &= 2 + a_1^4 + 2a_0a_2 + a_0^2a_2^2 + 2a_1^2a_2^2 + a_2^4 + 2a_1a_3 + 2a_0a_1a_2a_3 + a_3^2 + a_1^2a_3^2 \\ &\quad - 2a_1^2 - 2a_0a_1^2a_2 - a_2^2 - 2a_0a_2^3 - 2a_1^3a_3 - 2a_1a_2^2a_3. \end{aligned} \quad (3.8)$$

As a result, in  $\mathcal{M}(2)$  we now have

$$XY = a_1X + a_2Y \quad (3.9)$$

and

$$Y^2 = p_11 + p_2X + p_3Y + p_4X^2 \quad (3.10)$$

for suitable real scalars  $p_1, p_2, p_3, p_4$  (which depend upon  $a_0, a_1, a_2, a_3$ ); moreover,  $p_4 = -k$ . We will now build a flat moment matrix extension  $\mathcal{M}(3)$  of  $\mathcal{M}(2)$ . This will prove that  $\mathcal{M}(2)$  admits a 4-atomic representing measure; a fortiori,  $\beta^{(3)}$  also admits a 4-atomic representing measure, just as in Case 1 above.

To define  $\mathcal{M}(3)$ , we aim to preserve the property (RG). First, we observe that the columns  $1$ ,  $X$  and  $Y$  in  $\mathcal{M}(3)$  are obtained from those columns in  $\mathcal{M}(2)$  by adding suitable cubic and quartic moments. Moreover, the columns  $XY$  and



$Y^2$  are defined using (3.9) and (3.10), while the columns  $X^2Y$ ,  $XY^2$  and  $Y^3$  are obtained, via the functional calculus, from (3.9) and (3.10). For instance,

$$X^2Y := a_1X^2 + a_2XY = a_1X^2 + a_2(a_1X + a_2Y) = a_2a_1X + a_2^2Y + a_1X^2$$

and

$$\begin{aligned} Y^3 &:= p_1Y + p_2XY + p_3Y^2 + p_4X^2Y \\ &= p_1Y + p_2(a_1X + a_2Y) + p_3(p_11 + p_2X + p_3Y + p_4X^2) \\ &\quad + p_4(a_2a_1X + a_2^2Y + a_1X^2). \end{aligned}$$

It follows that both  $X^2Y$  and  $Y^3$  are linear combinations of the columns 1,  $X$ ,  $Y$  and  $X^2$ . Now, to define  $XY^2$  one can use either (3.9) or (3.10). However, property (RG) requires that both definitions of  $XY^2$  be compatible. In other words, the expressions

$$\begin{aligned} XY^2 &\equiv a_1XY + a_2Y^2 = a_2p_11 + (a_1^2 + a_2p_2)X + (a_1a_2 + a_2p_3)Y + a_2p_4X^2 \\ &\quad \text{(obtained using (3.9))} \end{aligned}$$

and

$$\begin{aligned} XY^2 &\equiv p_1X + p_2X^2 + p_3XY + p_4X^3 = (p_1 + a_1p_3)X + a_2p_3Y + p_2X^2 + p_4X^3 \\ &\quad \text{(obtained using (3.10))} \end{aligned}$$

must be identical. Since  $p_4 = -k \neq 0$ , we immediately get

$$X^3 = \frac{1}{p_4} [a_2p_11 + (a_1^2 + a_2p_2 - p_1 - a_1p_3)X + a_1a_2Y + (a_2p_4 - p_2)X^2], \quad (3.11)$$

which we can then use to define the column  $X^3$ . Close examination of (3.11) at the level of the fourth row in  $\mathcal{M}(3)$  leads to a formula for the quintic moment  $\beta_{50}$ . This value must then be inserted in the seventh row of  $X^2$ , to complete the definition of  $X^2$  in  $\mathcal{M}(3)$ . As a result, in the new moment matrix  $\mathcal{M}(3)$  we have exhibited each cubic column as a linear combination of columns associated with monomials of degree at most 2. This means that  $\mathcal{M}(3)$  is a flat extension of  $\mathcal{M}(2)$ , as desired. The proof is now complete.  $\square$

*Remark 3.3.* The quartic moment  $\beta_{04}$  defined by (3.8), is nonnegative, being a diagonal entry of the positive semidefinite matrix  $\mathcal{M}(2)$ . One can say more, however, by appealing to the theory of Semidefinite Programming (SDP). As is well known, a polynomial  $f \in \mathcal{P}_{2d}$  is a sum of squares if and only if  $f = \mathbf{z}^T Q \mathbf{z}$  for some square matrix  $Q \geq 0$ , where  $\mathbf{z}$  is the vector of monomials of degree less than or equal to  $d$ . If we let  $f \equiv f(a_0, a_1, a_2, a_3) := \beta_{04} - 1$  and  $\mathbf{y} := (1, a_2, a_3, a_1^2, a_2^2, a_0a_2, a_1a_3)$ , a calculation using *Mathematica* reveals that  $f \geq 0$  if and only if  $\mathbf{y}^T R \mathbf{y} \geq 0$ , where

$$R := \begin{pmatrix} 1 & 0 & 0 & -1 & -1 & 1 & 1 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 1 & 1 & -1 & -1 \\ -1 & 0 & 0 & 1 & 1 & -1 & -1 \\ 1 & 0 & 0 & -1 & -1 & 1 & 1 \\ 1 & 0 & 0 & -1 & -1 & 1 & 1 \end{pmatrix}. \quad (3.12)$$

Since  $R$  is a flat extension of its  $3 \times 3$  compression to the first three rows and columns, it is clear that  $R \geq 0$ . It follows that  $\beta_{04} \equiv f + 1 \geq 1 > 0$ .

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## REFERENCES

- [1] R. Curto and L. Fialkow, Solution of the truncated complex moment problem with flat data, *Memoirs Amer. Math. Soc.* no. 568, Amer. Math. Soc., Providence, 1996.
- [2] R. Curto and L. Fialkow, Flat extensions of positive moment matrices: Relations in analytic or conjugate terms, *Operator Th.: Adv. Appl.* 104(1998), 59–82.
- [3] R. Curto and L. Fialkow, Flat extensions of positive moment matrices: Recursively generated relations, *Memoirs Amer. Math. Soc.* no. 648, Amer. Math. Soc., Providence, 1998.
- [4] R. Curto and L. Fialkow, Solution of the singular quartic moment problem, *J. Operator Theory* 48 (2002), 315–354.
- [5] R. Curto and L. Fialkow, Solution of the truncated hyperbolic moment problem, *Integral Equations Operator Theory* 52 (2005), no. 2, 181–218.
- [6] R. Curto, L. Fialkow and H.M. Möller, The extremal truncated moment problem, *Integral Equations Operator Theory* 60 (2008), 177–200.
- [7] R. Curto and L. Fialkow, Recursively determined representing measures for bivariate truncated moment sequences, *J. Operator Theory* 70 (2013), no. 2, 401–436.
- [8] R. Curto, S. H. Lee, J. Yoon, A new approach to the 2-variable subnormal completion problem, *J. Math. Anal. Appl.* 370 (2010), no. 1, 270–283.
- [9] R. Curto and S. Yoo, Concrete solution of the nonsingular binary quartic moment problem, *to appear in Proc. Amer. Math. Soc.*
- [10] L. Fialkow and J. Nie, Positivity of Riesz functionals and solutions of quadratic and quartic moment problems, *J. Funct. Anal.* 258 (2010), 328–356.
- [11] D. P. Kimsey, *Matrix-Valued Moment Problems*, ProQuest LLC, Ann Arbor (2011). Ph.D. Thesis, Drexel University.
- [12] D. P. Kimsey, The cubic complex moment problem. *Integral Equations Operator Theory* 80 (2014), no. 3, 353–378.
- [13] D. P. Kimsey and H. J. Woerdeman, The truncated matrix-valued K-moment problem on  $\mathbb{R}^d$ ,  $\mathbb{C}^d$ , and  $\mathbb{T}^d$ , *Trans. Am. Math. Soc.* 365(10) (2013), 5393–5430.
- [14] J.L. Smul’jan, An operator Hellinger integral (Russian), *Mat. Sb.* 91 (1959), 381–430.
- [15] Wolfram Research, Inc., *Mathematica*, Version 9.0, Champaign, IL, 2013.

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