

Symmetry and Visual Hallucinations

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Klüver: We wish to stress . . . one point, namely, that under diverse conditions **the visual system** responds in terms of a limited number of **form constants**.

Planar Symmetry-Breaking

- **Euclidean symmetry**: translations, rotations, reflections
- **Symmetry-breaking** from translation invariant state in planar systems with **Euclidean symmetry** leads to
 - **Stripes**:
States invariant under translation in one direction
 - **Spots**:
States centered at lattice points

Sand Dunes in Namibian Desert



Mud Plains



Outline

1. Geometric Visual Hallucinations

2. Structure of Visual Cortex

Hubel and Wiesel hypercolumns; local and lateral connections; isotropy versus anisotropy

3. Pattern Formation in V1

Symmetry; Three models

4. Interpretation of Patterns in Retinal Coordinates

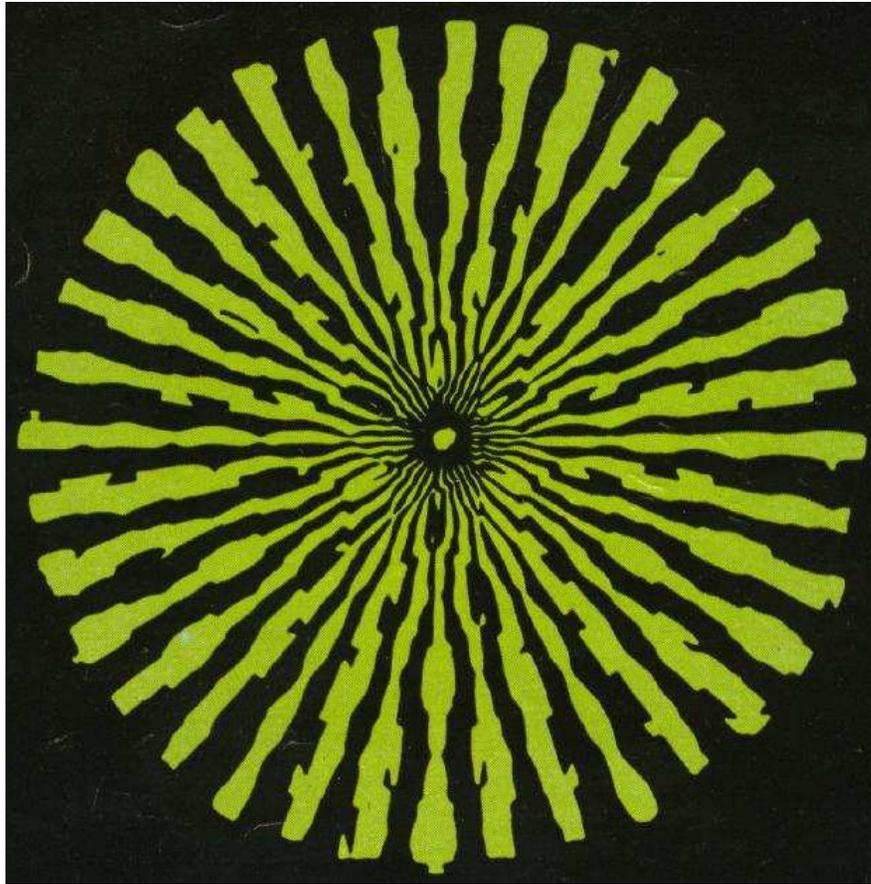
Visual Hallucinations

- Drug **uniformly** forces activation of cortical cells
- Leads to **spontaneous** pattern formation on cortex
- Map from V1 to retina;
translates pattern on cortex to visual image

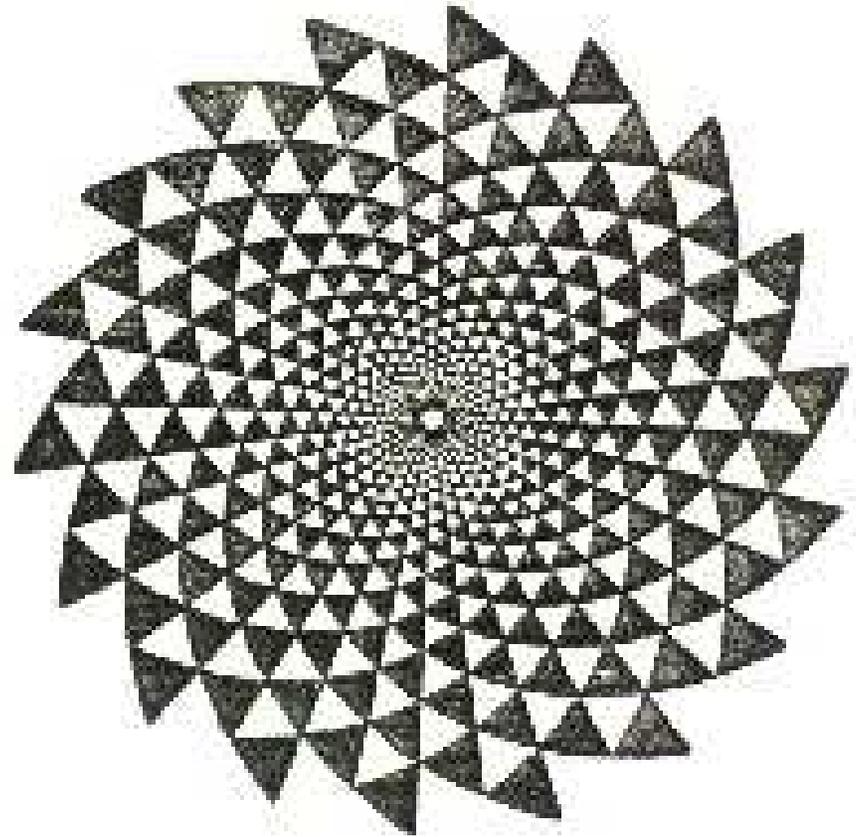
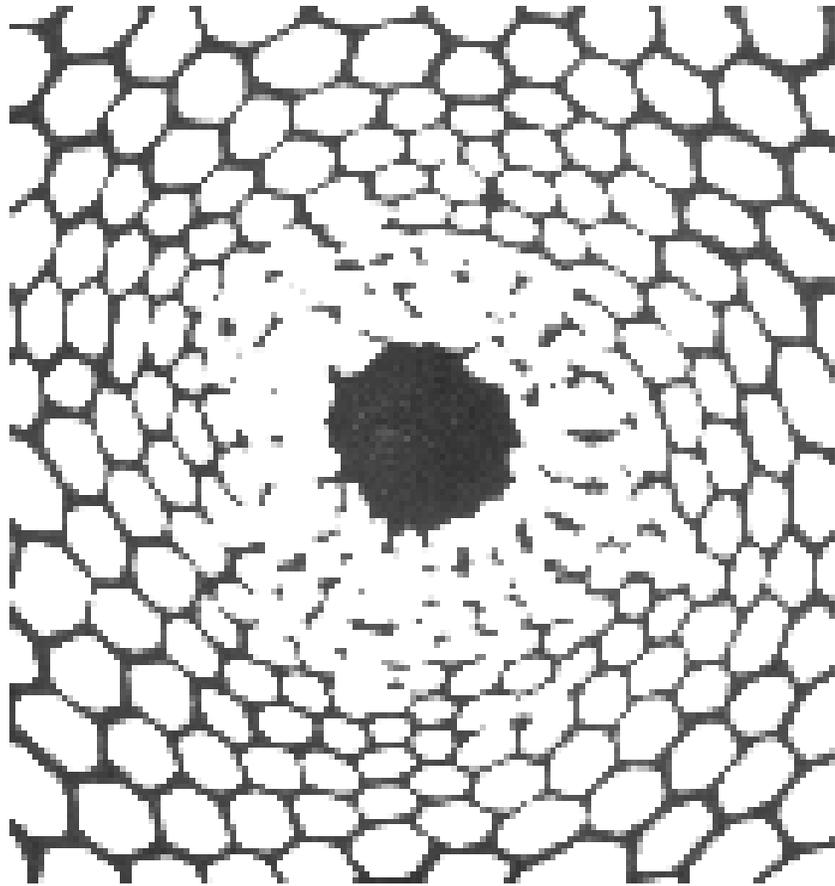
Visual Hallucinations

- Drug **uniformly** forces activation of cortical cells
- Leads to **spontaneous** pattern formation on cortex
- Map from V1 to retina;
translates pattern on cortex to visual image
- Patterns fall into four ***form constants*** (Klüver, 1928)
 - **tunnels** and **funnels**
 - **spirals**
 - **lattices** includes **honeycombs** and **phosphenes**
 - **cobwebs**

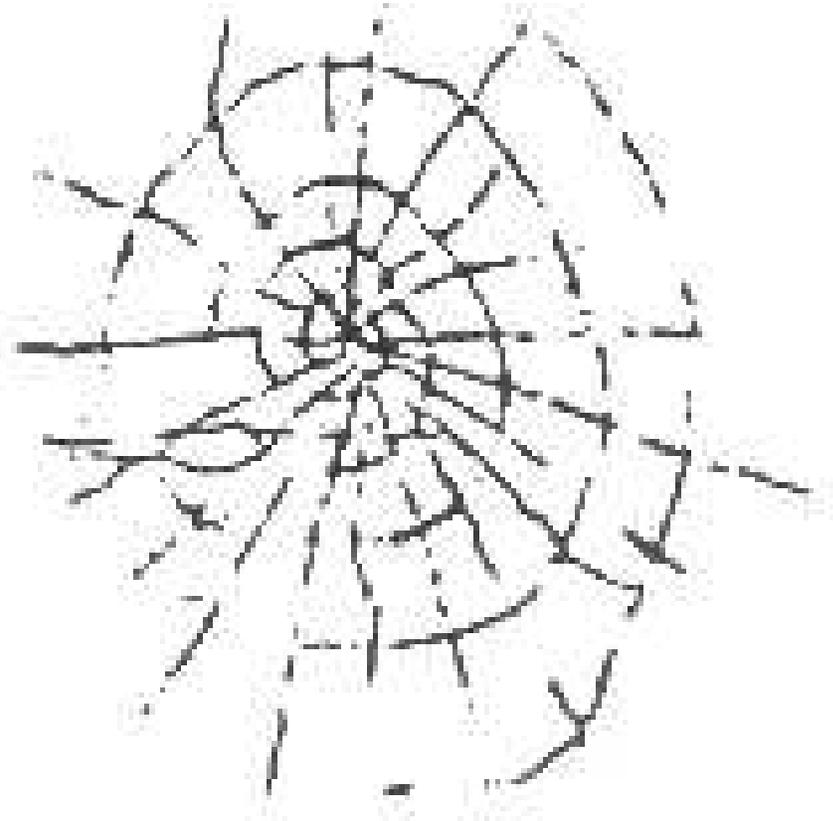
Funnel and Spirals



Lattices: Honeycombs & Phosphenes

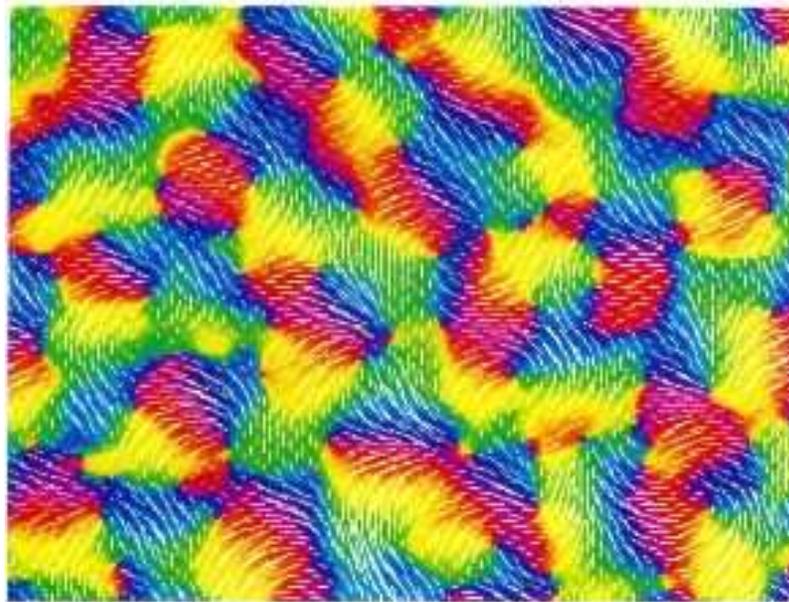


Cobwebs



Orientation Sensitivity of Cells in V1

- Most V1 cells sensitive to *orientation* of contrast edge

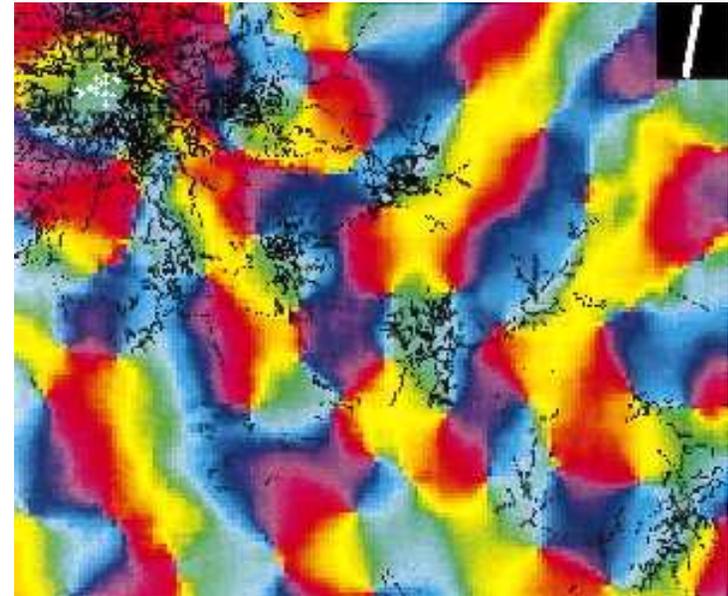
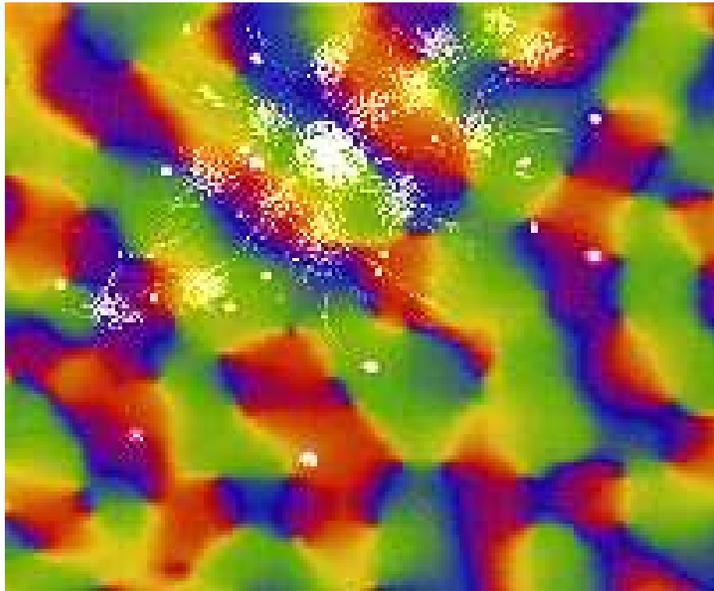


Distribution of orientation preferences in Macaque V1 (Blasdel)

- Hubel and Wiesel, 1974
Each millimeter there is a *hypercolumn* consisting of orientation sensitive cells in every direction preference

Structure of Primary Visual Cortex (V1)

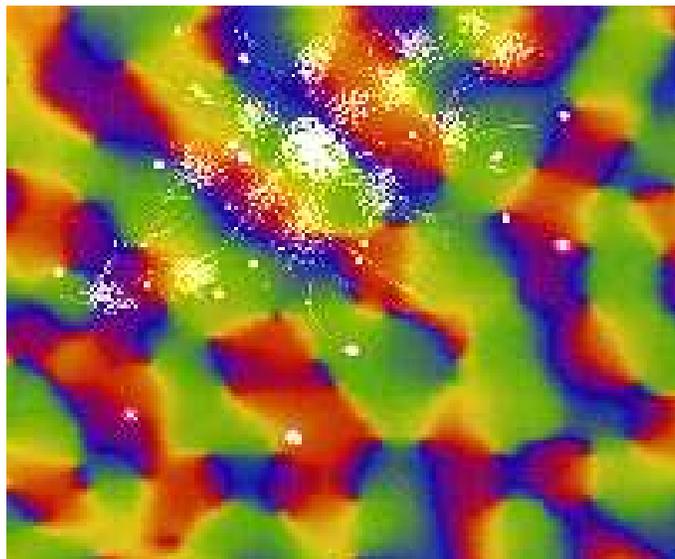
- Optical imaging exhibits pattern of connection



V1 lateral connections: Macaque (left, Blasdel) and Tree Shrew (right, Fitzpatrick)

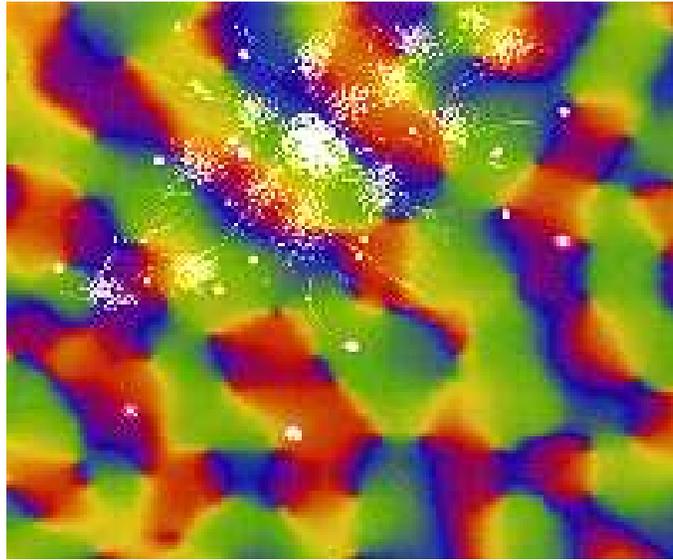
- Two kinds of coupling: **local** and **lateral**
 - local**: cells $< 1mm$ connect with most neighbors
 - lateral**: cells make contact each mm along axons; connections in direction of cell's preference

Anisotropy in Lateral Coupling

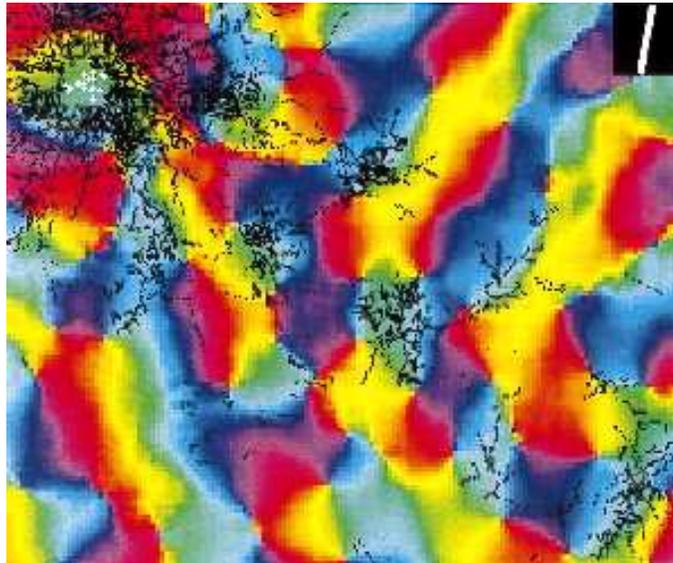


- Macaque: most anisotropy due to stretching in direction orthogonal to ocular dominance columns. Anisotropy is weak.

Anisotropy in Lateral Coupling

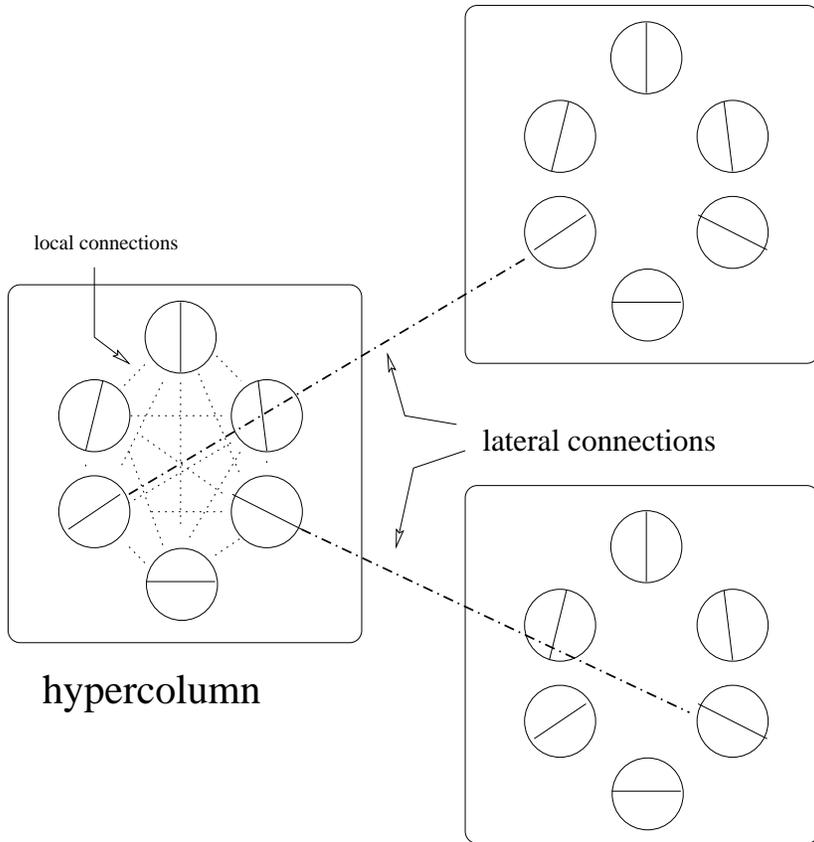


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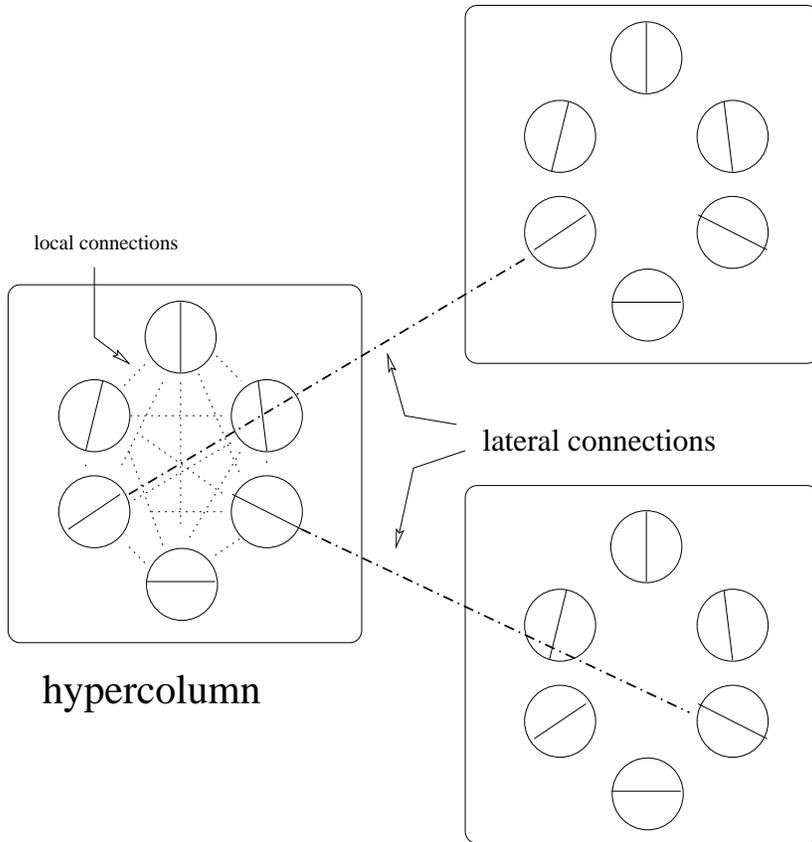


- Tree shrew: anisotropy pronounced

Action of Euclidean Group: Anisotropy



Action of Euclidean Group: Anisotropy



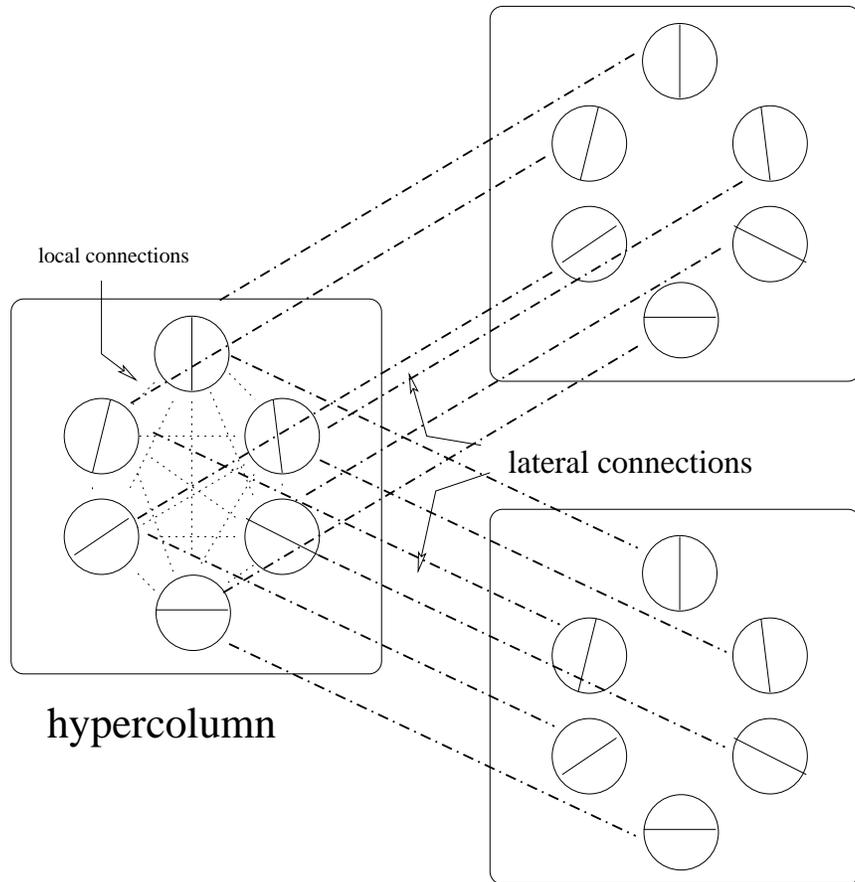
- Abstract **physical space** of V1 is $\mathbb{R}^2 \times S^1$ — not \mathbb{R}^2
Hypercolumn becomes circle of orientations
- Euclidean group on \mathbb{R}^2 :
translations, rotations, reflections
- Euclidean groups acts on $\mathbb{R}^2 \times S^1$ by

$$T_y(x, \varphi) = (T_y x, \varphi)$$

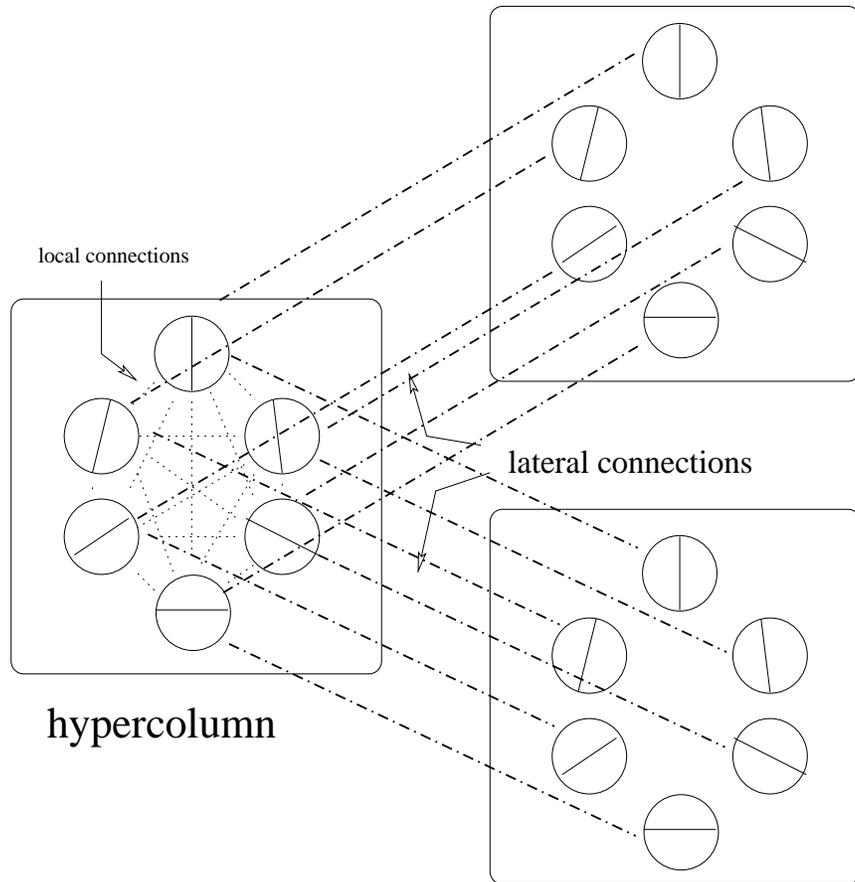
$$R_\theta(x, \varphi) = (R_\theta x, \varphi + \theta)$$

$$\kappa(x, \varphi) = (\kappa x, -\varphi)$$

Isotropic Lateral Connections



Isotropic Lateral Connections



- New $O(2)$ symmetry

$$\hat{\phi}(x, \varphi) = (x, \varphi + \hat{\phi})$$

- Weak anisotropy is forced symmetry breaking of

$$\mathbf{E}(2) \dot{+} \mathbf{O}(2) \rightarrow \mathbf{E}(2)$$

Three Models

- $E(2)$ acting on \mathbb{R}^2 (Ermentrout-Cowan)
neurons located at each point x

Activity variable: $a(x)$ = voltage potential of neuron

Pattern given by threshold $a(x) > v_0$

Three Models

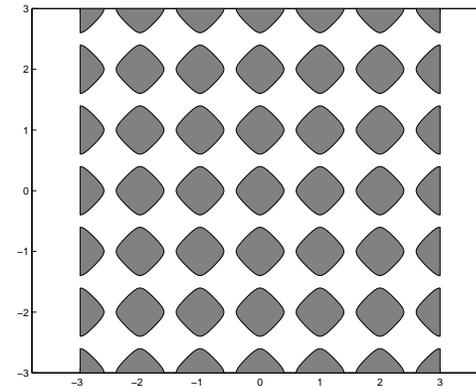
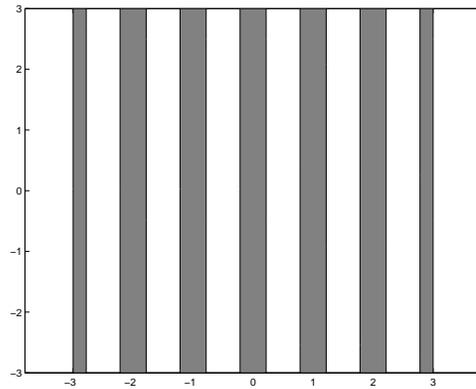
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hypercolumns located at x ; neurons tuned to φ
strongly anisotropic lateral connections
Activity variable: $a(x, \varphi)$
Pattern given by winner-take-all

Three Models

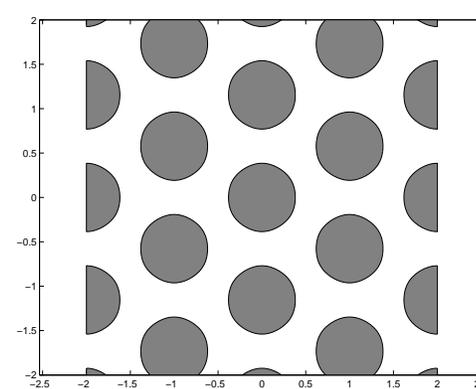
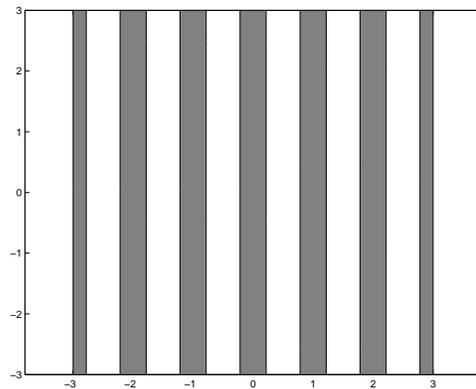
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Pattern given by winner-take-all
- Symmetry breaking: $E(2) \dot{+} O(2) \rightarrow E(2)$
weakly anisotropic lateral coupling
Activity variable: $a(x, \varphi)$
Pattern given by winner-take-all

Planforms For Ermentrout-Cowan

Threshold Patterns



Square lattice: stripes and squares



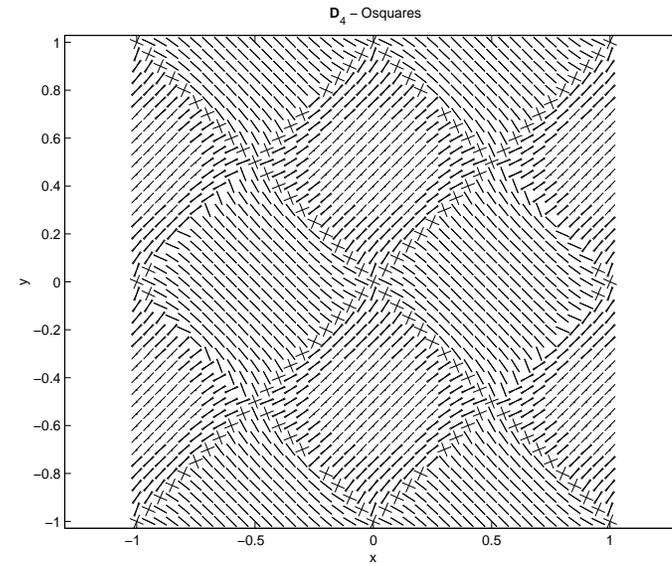
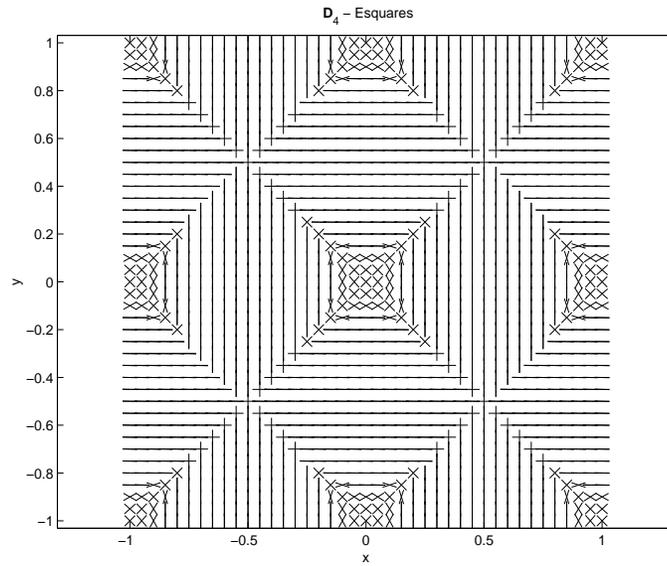
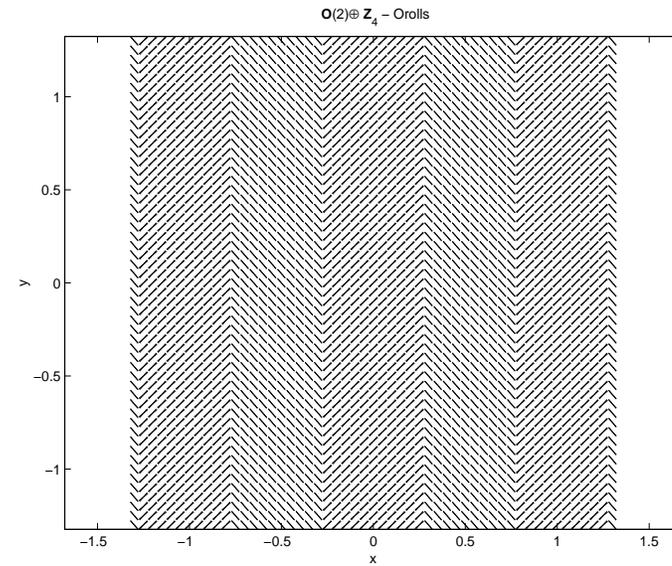
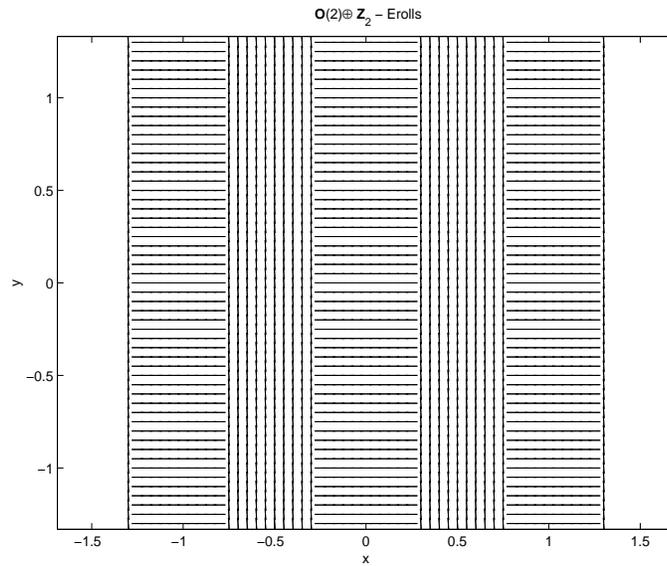
Hexagonal lattice: stripes and hexagons

Winner-Take-All Strategy

Creation of Line Fields

- **Given:** Activity $a(\mathbf{x}, \varphi)$ of neuron in hypercolumn at \mathbf{x} sensitive to direction φ
- **Assumption:** Most active neuron in hypercolumn suppresses other neurons in hypercolumn
- **Consequence:** For all \mathbf{x} find direction $\varphi_{\mathbf{x}}$ where activity is maximum
- **Planform:** Line segment at each \mathbf{x} oriented at angle $\varphi_{\mathbf{x}}$

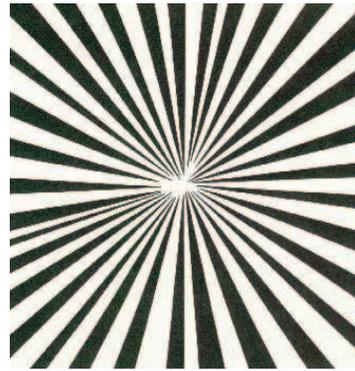
Planforms For Bressloff-Cowan



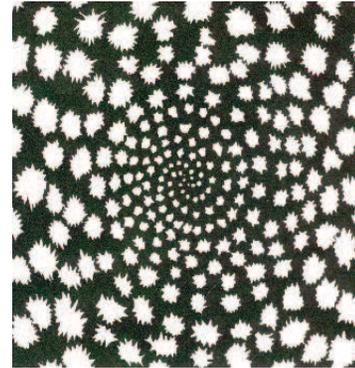
Cortex to Retina

- Neurons on cortex are **uniformly** distributed
- Neurons in retina fall off by $1/r^2$ from fovea
- Unique **angle preserving** map takes uniform density square to $1/r^2$ density disk: **complex exponential**
- Straight lines on cortex \mapsto **circles, logarithmic spirals, and rays** in retina

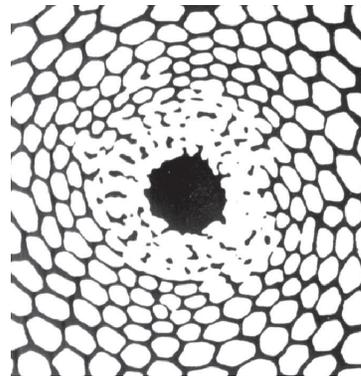
Visual Hallucinations



(I)



(II)



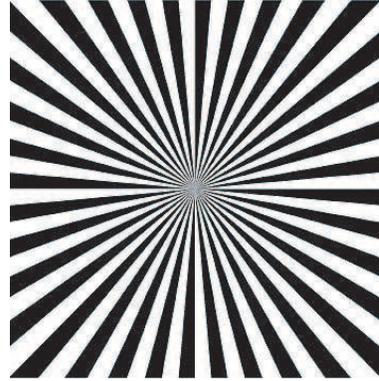
(III)



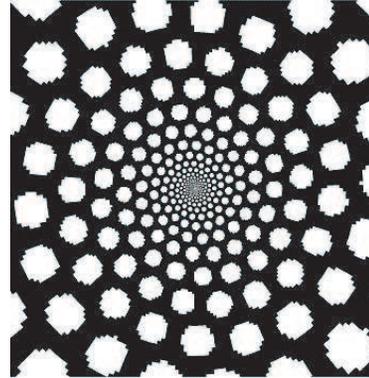
(IV)

(I) **funnel** and (II) **spiral** images LSD [Siegel & Jarvik, 1975], (III) **honeycomb** marijuana [Clottes & Lewis-Williams (1998)], (IV) **cobweb** petroglyph [Patterson, 1992]

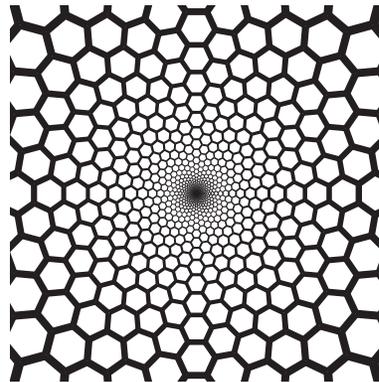
Planforms in the Visual Field



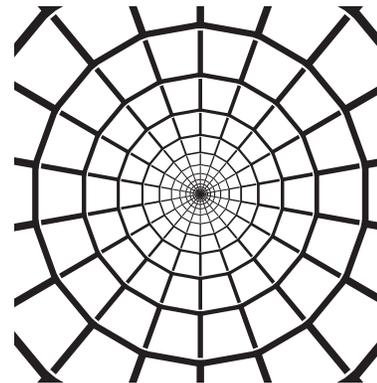
(a)



(b)



(c)



(d)

Visual field planforms

Weakly Anisotropic Coupling

- In addition to equilibria found in Bressloff-Cowan model there exist **periodic solutions** that emanate from steady-state bifurcation

1. **Rotating** Spirals
2. **Tunneling** Blobs **Tunneling Spiraling** Blobs
3. **Pulsating** Blobs

Pattern Formation Outline

1. **Bifurcation Theory with Symmetry**
 - Equivariant Branching Lemma
 - Model independent analysis
2. **Translations** lead to plane waves
3. **Planforms**: Computation of eigenfunctions

Primer on Steady-State Bifurcation

- Solve $\dot{x} = f(x, \lambda) = 0$ where $f : \mathbf{R}^n \times \mathbf{R} \rightarrow \mathbf{R}^n$
- Local theory: Assume $f(0, 0) = 0$ & find solns near $(0, 0)$
- If $L = (d_x f)_{0,0}$ nonsingular, IFT implies unique soln $x(\lambda)$

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- If $L = (d_x f)_{0,0}$ nonsingular, IFT implies unique soln $x(\lambda)$
- Bifurcation of steady states $\iff \ker L \neq \{0\}$
- Reduction theory implies that steady-states are found by solving $\varphi(y, \lambda) = 0$ where

$$\varphi : \ker L \times \mathbf{R} \rightarrow \ker L$$

Equivariant Steady-State Bifurcation

Let $\gamma : \mathbf{R}^n \rightarrow \mathbf{R}^n$ be linear

- γ is a symmetry iff $\gamma(\text{soln}) = \text{soln}$ iff $f(\gamma x, \lambda) = \gamma f(x, \lambda)$
- Chain rule $\implies L\gamma = \gamma L \implies \ker L$ is γ -invariant
- **Theorem:** Fix symmetry group Γ . Generically $\ker L$ is an absolutely irreducible representation of Γ
- Reduction implies that there is a unique steady-state bifurcation theory for each absolutely irreducible rep

Equivariant Bifurcation Theory

- Let $\Sigma \subset \Gamma$ be a subgroup
- $\text{Fix}(\Sigma) = \{x \in \ker L : \sigma x = x \quad \forall \sigma \in \Sigma\}$
- Σ is *axial* if $\dim \text{Fix}(\Sigma) = 1$

- Equivariant Branching Lemma:

Generically, there exists a branch of solutions with Σ symmetry for every axial subgroup Σ

- **MODEL INDEPENDENT**

Solution types do not depend on the equation — only on the symmetry group and its representation on $\ker L$

Translations

- Let $W_{\mathbf{k}} = \{u(\varphi)e^{i\mathbf{k}\cdot\mathbf{x}} + \text{c.c.}\}$ $\mathbf{k} \in \mathbf{R}^2 =$ **wave vector**

- **Translations** act on $W_{\mathbf{k}}$ by

$$T_{\mathbf{y}}(u(\varphi)e^{i\mathbf{k}\cdot\mathbf{x}}) = u(\varphi)e^{i\mathbf{k}\cdot(\mathbf{x}+\mathbf{y})} = \left[e^{i\mathbf{k}\cdot\mathbf{y}} u(\varphi) \right] e^{i\mathbf{k}\cdot\mathbf{x}}$$

- $L : W_{\mathbf{k}} \rightarrow W_{\mathbf{k}}$

Eigenfunctions of L have *plane wave factors*

Reflections

- Choose **REFLECTION** ρ so that $\rho\mathbf{k} = \mathbf{k}$

$$\rho \left(u(\varphi) e^{i\mathbf{k}\cdot\mathbf{x}} \right) = \rho(u(\varphi)) e^{i\mathbf{k}\cdot\mathbf{x}}$$

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- $\rho^2 = 1$ implies

$$W_{\mathbf{k}} = W_{\mathbf{k}}^+ \oplus W_{\mathbf{k}}^-$$

where ρ acts as $+1$ on $W_{\mathbf{k}}^+$ and -1 on $W_{\mathbf{k}}^-$

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- Eigenfunctions are **even** or **odd**. When $\mathbf{k} = (1, 0)$

$$u(-\varphi) = u(\varphi) \quad u \in W_{\mathbf{k}}^+$$

$$u(-\varphi) = -u(\varphi) \quad u \in W_{\mathbf{k}}^-$$

Rotations

- **Rotations** act on spaces $W_{\mathbf{k}}$

$$R_{\theta} \left(u(\varphi) e^{i\mathbf{k}\cdot\mathbf{x}} \right) = R_{\theta}(u(\varphi)) e^{iR_{\theta}(\mathbf{k})\cdot\mathbf{x}}$$

Therefore

$$R_{\theta}(W_{\mathbf{k}}) = W_{R_{\theta}(\mathbf{k})}$$

Therefore $\ker L$ is **∞ -dimensional**

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- **Double-periodicity**: Look for solutions on planar lattice

$$\mathcal{F}_{\mathcal{L}} = \{f \in \mathcal{F} : f(\mathbf{x} + \ell) = f(\mathbf{x}) \quad \forall \ell \in \mathcal{L}\}$$

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- **Finite number of rotations**: $\ker L$ is **finite**-dimensional
- Choose lattice size so **shortest** dual vectors are **critical**

Axials in Ermentrout-Cowan Model

Name	Planform Eigenfunction
stripes	$\cos x$
squares	$\cos x + \cos y$
hexagons	$\cos(\mathbf{k}_0 \cdot \mathbf{x}) + \cos(\mathbf{k}_1 \cdot \mathbf{x}) + \cos(\mathbf{k}_2 \cdot \mathbf{x})$

$$\mathbf{k}_0 = (1, 0)$$

$$\mathbf{k}_1 = \frac{1}{2}(-1, \sqrt{3})$$

$$\mathbf{k}_2 = \frac{1}{2}(-1, -\sqrt{3})$$

Axials in Bressloff-Cowan Model

Name	Planform Eigenfunction	u
squares	$u(\varphi) \cos x + u\left(\varphi - \frac{\pi}{2}\right) \cos y$	even
stripes	$u(\varphi) \cos x$	even
hexagons	$\sum_{j=0}^2 u\left(\varphi - j\pi/3\right) \cos(\mathbf{k}_j \cdot \mathbf{x})$	even
square stripes	$u(\varphi) \cos x - u\left(\varphi - \frac{\pi}{2}\right) \cos y$ $u(\varphi) \cos x$	odd odd
hexagons	$\sum_{j=0}^2 u\left(\varphi - j\pi/3\right) \cos(\mathbf{k}_j \cdot \mathbf{x})$	odd
triangles	$\sum_{j=0}^2 u\left(\varphi - j\pi/3\right) \sin(\mathbf{k}_j \cdot \mathbf{x})$	odd
rectangles	$u\left(\varphi - \frac{\pi}{3}\right) \cos(\mathbf{k}_1 \cdot \mathbf{x}) - u\left(\varphi + \frac{\pi}{3}\right) \cos(\mathbf{k}_2 \cdot \mathbf{x})$	odd

How to Find Amplitude Function $u(\varphi)$

- **Isotropic connections** imply EXTRA $O(2)$ symmetry
- $O(2)$ decomposes $W_{\mathbf{k}}$ into sum of irreducible subspaces

$$W_{\mathbf{k},p} = \{ze^{p\varphi i}e^{i\mathbf{k}\cdot\mathbf{x}} + \text{c.c.} : z \in \mathbf{C}\} \cong \mathbf{R}^2$$

Eigenfunctions lie in $W_{\mathbf{k},p}$ for some p

- $W_{\mathbf{k},p}^+ = \{\cos(p\varphi)e^{i\mathbf{k}\cdot\mathbf{x}}\}$ **even case**
- $W_{\mathbf{k},p}^- = \{\sin(p\varphi)e^{i\mathbf{k}\cdot\mathbf{x}}\}$ **odd case**

- With **weak anisotropy**

$$u(\varphi) \approx \cos(p\varphi) \quad \text{or} \quad u(\varphi) \approx \sin(p\varphi)$$

Rotating waves

- Suppose $\text{Fix}(\Sigma)$ is two-dimensional
Suppose $N_{\Gamma}(\Sigma) = \Sigma \times \mathbf{SO}(2)$
- Then generically solutions are rotating waves of a pattern with Σ symmetry
- Leads to **rotating spirals** and **tunnels**
- Suppose $N_{\Gamma}(\Sigma) = \Sigma \times \mathbf{D}_4$
- Leads to **pulsating** solutions