

## POISSON'S EQUATION - DISCRETIZATION

The Dirichlet boundary value problem for Poisson's equation is given by

$$\begin{aligned}\Delta u(x, y) &= g(x, y), & (x, y) \in R \\ u(x, y) &= f(x, y), & (x, y) \in \Gamma\end{aligned}\quad (1)$$

where  $R$  is a planar region and  $\Gamma$  is the boundary of  $R$ . In the book in §8.8, we examine this problem for  $R = \{(x, y) \mid 0 < x, y < 1\}$ .

For an integer  $N > 1$ , we introduce a mesh size  $h = 1/N$ . With it, we define a rectangular mesh on  $\bar{R} = R \cup \Gamma$ :

$$(x_j, y_k) = \left(\frac{j}{N}, \frac{k}{N}\right), \quad j, k = 0, 1, \dots, N$$

At mesh points (=grid points) inside of  $R$ , we approximate the equation  $\Delta u = g$ . To do so, we recall the approximation

$$G''(x) = \frac{G(x+h) - 2G(x) + G(x-h)}{h^2} - \frac{h^2}{12}G^{(4)}(\xi)$$

with some  $\xi \in [x-h, x+h]$ . This assumes  $G$  is four times continuously differentiable on  $[x-h, x+h]$ .

Since

$$\Delta u(x, y) = \frac{\partial^2 u(x, y)}{\partial x^2} + \frac{\partial^2 u(x, y)}{\partial y^2}$$

we obtain at each  $(x_j, y_k) \in R$  that

$$\begin{aligned} & \frac{u(x_{j+1}, y_k) - 2u(x_j, y_k) + u(x_{j-1}, y_k)}{h^2} \\ & + \frac{u(x_j, y_{k+1}) - 2u(x_j, y_k) + u(x_j, y_{k-1})}{h^2} \\ & - \frac{h^2}{12} \left[ \frac{\partial^4 u(\xi_j, y_k)}{\partial x^4} + \frac{\partial^4 u(x_j, \eta_k)}{\partial y^4} \right] \\ & = g(x_j, y_k) \end{aligned} \quad (2)$$

with  $\xi_j \in [x_{j-1}, x_{j+1}]$ ,  $\eta_k \in [y_{k-1}, y_{k+1}]$ . We can delete the error terms involving the fourth derivatives of  $u$ . This leads to a new set of equations for approximating unknown  $u_h \approx u$ :

$$\{u_h(x_j, y_k) \mid 1 < j, k < N\}$$

The values of  $u_h$  at boundary mesh points on  $\Gamma$  are given by

$$u_h(x_j, y_k) = f(x_j, y_k), \quad (x_j, y_k) \in \Gamma \quad (3)$$

This leads to the linear system

$$\begin{aligned} & \frac{u_h(x_{j+1}, y_k) - 2u_h(x_j, y_k) + u_h(x_{j-1}, y_k)}{h^2} \\ & + \frac{u_h(x_j, y_{k+1}) - 2u_h(x_j, y_k) + u_h(x_j, y_{k-1})}{h^2} \\ & = g(x_j, y_k), \quad (x_j, y_k) \in R \end{aligned}$$

$$u_h(x_j, y_k) = f(x_j, y_k), \quad (x_j, y_k) \in \Gamma \quad (4)$$

For interior mesh points, we can simplify this to

$$\begin{aligned} & u_h(x_j, y_{k-1}) + u_h(x_{j-1}, y_k) - 4u_h(x_j, y_k) \\ & + u_h(x_{j+1}, y_k) + u_h(x_j, y_{k+1}) = h^2 g(x_j, y_k) \\ & \quad \quad \quad (x_j, y_k) \in R \end{aligned} \quad (5)$$

Thus we have a system of  $(N + 1)^2$  linear equations in the  $(N + 1)^2$  unknowns

$$\{u_h(x_j, y_k) \mid 1 \leq j, k \leq N\}$$

If  $u(x, y)$  is four times continuously differentiable over  $\overline{R}$ , it can be shown that for some  $c$ ,

$$\max_R |u(x_j, y_k) - u_h(x_j, y_k)| \leq c h^2$$

## JACOBI ITERATION

For  $m = 0, 1, 2, \dots$ , define

$$u_h^{(m+1)}(x_j, y_k) = \frac{1}{4} \left\{ u_h^{(m)}(x_j, y_{k-1}) + u_h^{(m)}(x_{j-1}, y_k) \right. \\ \left. + u_h^{(m)}(x_{j+1}, y_k) + u_h^{(m)}(x_j, y_{k+1}) - h^2 g(x_j, y_k) \right\} \\ (x_j, y_k) \in R \quad (6)$$

For values of  $u_h^{(m)}$  at boundary points, use the given boundary conditions as in (3).

We can write the iteration in matrix-vector form as

$$\mathbf{u}_h^{(m+1)} = Q \mathbf{u}_h^{(m)} + \mathbf{G}_h$$

$$Q = \frac{1}{4} \begin{bmatrix} T & I & 0 & \cdots & 0 & 0 \\ I & T & I & & & \\ 0 & I & T & I & & \vdots \\ \vdots & & & \ddots & & \\ 0 & & & I & T & I \\ 0 & \cdots & & & I & T \end{bmatrix}$$

with  $I$  the identity of order  $N - 1$ ,

$$T = \begin{bmatrix} 0 & 1 & \cdots & 0 \\ 1 & 0 & 1 & \\ \vdots & & \ddots & \vdots \\ & & 1 & 0 & 1 \\ 0 & \cdots & 1 & 0 \end{bmatrix}$$

$$\mathbf{u}_h^{(m)} = \begin{bmatrix} \mathbf{u}_h^{(m)}(*, y_1) \\ \mathbf{u}_h^{(m)}(*, y_2) \\ \vdots \\ \mathbf{u}_h^{(m)}(*, y_{N-1}) \end{bmatrix}$$

$$\mathbf{u}_h^{(m)}(*, y_k) = \begin{bmatrix} u_h^{(m)}(x_1, y_k) \\ u_h^{(m)}(x_2, y_k) \\ \vdots \\ u_h^{(m)}(x_{N-1}, y_k) \end{bmatrix}$$

$$\mathbf{G}_h = \begin{bmatrix} \frac{1}{4} \left\{ \mathbf{B}(y_1) - h^2 \mathbf{G}_h(*, y_1) \right\} \\ \frac{1}{4} \left\{ \mathbf{B}(y_2) - h^2 \mathbf{G}_h(*, y_2) \right\} \\ \vdots \\ \frac{1}{4} \left\{ \mathbf{B}(y_{N-1}) - h^2 \mathbf{G}_h(*, y_{N-1}) \right\} \end{bmatrix}$$

$$\mathbf{B}(y_k) = \begin{bmatrix} f(x_0, y_k) \\ 0 \\ \vdots \\ 0 \\ f(x_N, y_k) \end{bmatrix}$$

$$\mathbf{G}_h(*, y_k) = \begin{bmatrix} g(x_1, y_k) \\ g(x_2, y_k) \\ \vdots \\ g(x_{N-1}, y_k) \end{bmatrix}$$

The original discretization of (4)-(5) can be rewritten as

$$\mathbf{u}_h = \mathbf{Q}\mathbf{u}_h + \mathbf{G}_h$$

with  $\mathbf{u}_h$  written in accordance with the above definition of  $\mathbf{u}_h^{(m)}$ .

$$\mathbf{u}_h = Q\mathbf{u}_h + \mathbf{G}_h \quad (7)$$

$$\mathbf{u}_h^{(m+1)} = Q\mathbf{u}_h^{(m)} + \mathbf{G}_h \quad (8)$$

In this block matrix description of the system being solved, each block  $\mathbf{u}_h(*, y_k)$  corresponds to the horizontal row of interior grid points at  $y = y_k$ , with the entries in  $\mathbf{u}_h(*, y_k)$  being ordered in correspondence to the grid points being numbered from left to right in the grid. Note that  $Q$  is a symmetric matrix of order  $(N - 1)^2$ .

If we write our linear system (7) in the form  $Ax = b$ , we have  $A = I - Q$ , again a symmetric matrix. With reference to the Gauss-Seidel iteration, note that the diagonal elements of  $A$  are positive.

For the iteration error  $e_h^{(m)}$  in the Jacobi iteration, subtract (8) from (7), obtaining

$$e_h^{(m+1)} = Qe_h^{(m)}, \quad m = 0, 1, \dots \quad (9)$$

The matrix  $Q$  is the earlier matrix  $M$  of our general framework in the case of the Gauss-Jacobi iteration. Thus we have convergence of the Jacobi iteration if and only if

$$r_\sigma(Q) < 1$$

Note in this case that

$$\|Q\|_1 = \|Q\|_\infty = 1$$

Thus we are driven to looking at  $r_\sigma(Q)$ , and in this case  $\|Q\|_2 = r_\sigma(Q)$ .



## GAUSS-SEIDEL ITERATION

For  $m = 0, 1, 2, \dots$ , define

$$\begin{aligned} u_h^{(m+1)}(x_j, y_k) = & \frac{1}{4} \left\{ u_h^{(m+1)}(x_j, y_{k-1}) \right. \\ & + u_h^{(m+1)}(x_{j-1}, y_k) + u_h^{(m)}(x_{j+1}, y_k) \\ & \left. + u_h^{(m)}(x_j, y_{k+1}) \right\} - h^2 g(x_j, y_k), \quad (x_j, y_k) \in R \end{aligned} \tag{10}$$

For values of  $u_h^{(m)}$  at boundary points, use the given boundary conditions as in (3).

We can use the above matrices to also write this iteration in matrix-vector format. In the language of the general framework given earlier, it amounts to letting  $N$  be the lower triangular portion of  $A = I - Q$  and  $P = N - A$ . Then write

$$N\mathbf{u}_h^{(m+1)} = \mathbf{G}_h + P\mathbf{u}_h^{(m)}$$

## RATES OF CONVERGENCE

Introduce

$$\xi = 1 - 2 \sin^2 \left( \frac{\pi}{2N} \right) \quad (11)$$

$$\omega^* = \frac{2}{1 + \text{sqrt} (1 - \xi^2)}$$

Then for the rates of convergence of the various iteration methods we have studied when applied to solving (7),  $\mathbf{u}_h = Q\mathbf{u}_h + \mathbf{G}_h$ , we have

Rate of Convergence	
Method	$r_\sigma(M)$
Jacobi	$\xi$
Gauss-Seidel	$\xi^2$
SOR	$\omega^* - 1$

## THE RATE OF CONVERGENCE OF SOR

In the language of the general framework given earlier, the error equation for the SOR method will be

$$e_h^{(m+1)} = M(\omega) e_h^{(m)}, \quad m = 0, 1, \dots \quad (12)$$

The eigenvalues of  $M(\omega)$  can be computed explicitly in this case, yielding

$$r_\sigma(M(\omega)) = \begin{cases} \frac{\omega\xi}{2} + \text{sqrt} \left( \left( \frac{\omega\xi}{2} \right)^2 + 4(1-\omega) \right), & 0 \leq \omega \leq \omega^* \\ \omega - 1, & \omega^* \leq \omega \leq 2 \end{cases}$$

We give graphs for a few cases of the discretization parameter  $N$ .

## COSTS

Recall the notation from §8.7 regarding costs:

$$\text{Cost}(c, \varepsilon, n) = m^* \nu(n)$$

with  $c$  the factor by which the error was decreasing per iterate,  $n$  the order of the system,  $\varepsilon$  the error tolerance in

$$\|x - x^{(m)}\| \leq \varepsilon \|x - x^{(0)}\|$$

$\nu(n)$  the number of arithmetic operations per iterate calculation, and

$$m^* = \frac{-\log \varepsilon}{R(c)}, \quad R(c) = -\log c$$

For  $c = 1 - \delta$  with  $\delta$  small, we have

$$R(c) \doteq \delta \quad c = 1 - \delta, \quad \delta \approx 0$$

Note that for Poisson example,  $\nu(n) \approx 5n$ , and thus the total cost of attaining the desired accuracy is

$$\text{Cost}(c, \varepsilon, n) = 5nm^*$$

## APPLICATION TO SOLVING POISSON'S EQUATION

Gauss-Jacobi :

$$c = \xi = 1 - \frac{1}{2}\pi^2 h^2 + O(h^4)$$
$$\delta \approx \frac{1}{2}\pi^2 h^2$$

Gauss-Seidel :

$$c = \xi^2 = 1 - \pi^2 h^2 + O(h^4)$$
$$\delta \approx \pi^2 h^2$$

SOR :

$$c = \omega^* - 1 = 1 - 2\pi h + O(h^2)$$
$$\delta \approx 2\pi h$$

The improvement in the needed numbers of iterates is as follows, along with total costs for a desired accuracy.

Gauss-Jacobi vs. Gauss-Seidel : Gauss-Seidel requires approximately half the iterates required of Gauss-Jacobi. For Gauss-Seidel iteration,

$$\begin{aligned} \text{Cost}(c, \varepsilon, n) &\approx -\frac{5n \log \varepsilon}{\pi^2 h^2} = -\frac{5(N-1)^2 \log \varepsilon}{\pi^2 (N-1)^{-2}} \\ &\approx -\frac{5N^4 \log \varepsilon}{\pi^2} \end{aligned}$$

Gauss-Seidel vs. SOR : The factor giving the decrease in the number of iterates to be computed is

$$\frac{2\pi h}{\pi^2 h^2} = \frac{2}{\pi h}$$

Thus you are increasingly better off as  $h \rightarrow 0$ . For the cost of SOR,

$$\begin{aligned} \text{Cost}(c, \varepsilon, n) &\approx -\frac{5n \log \varepsilon}{2\pi h} = -\frac{5(N-1)^2 \log \varepsilon}{2\pi (N-1)^{-1}} \\ &\approx -\frac{5N^3 \log \varepsilon}{2\pi} \end{aligned}$$

The costs increase rapidly with  $N$ .

Let  $K_J$ ,  $K_{GS}$ ,  $K_{SOR}$  denote the number of iterations needed to reduce the iteration error by a factor of 10, for the Jacobi, Gauss-Seidel, and SOR methods, respectively. The following table shows the results for several values of  $N$ .

	$N = 10$	$N = 20$	$N = 40$	$N = 80$
$\xi$	.9510565	.9876883	.9969173	.99922904
$\xi^2$	.9045085	.9755283	.9938442	.99845867
$\omega^* - 1$	.5278640	.7294538	.8544978	.92444658
$K_J$	46	186	746	2985
$K_{GS}$	23	93	373	1493
$K_{SOR}$	4	8	15	31

## AN INITIAL GUESS

How to generate an initial guess  $u_h^{(0)}$ ? Since we know the solution on the boundary,

$$u_h(x_j, y_k) = f(x_j, y_k), \quad (x_j, y_k) \in \Gamma$$

We use an “interpolant” of this data to extend the boundary values to all of  $R$ . Define

$$\begin{aligned} u_h^{(0)}(x, y) = & (1 - x)f(0, y) + xf(1, y) \\ & + (1 - y)f(x, 0) + yf(x, 1) \\ & - [(1 - y)(1 - x)f(0, 0) + xyf(1, 1) \\ & + (1 - y)xf(1, 0) + y(1 - x)f(0, 1)] \end{aligned} \tag{13}$$



## LINE ITERATION

Recall the discretization of (8.8.5):

$$u_h(x_j, y_k) = \frac{1}{4} \left\{ u_h(x_{j+1}, y_k) + u_h(x_j, y_{k+1}) \right. \\ \left. + u_h(x_{j-1}, y_k) + u_h(x_j, y_{k-1}) \right\} \\ - \frac{h^2}{4} g(x_j, y_k), \quad 1 \leq j, k \leq N - 1$$

Previously we solved individually for a new value of  $u_h(x_j, y_k)$ . Consider solving simultaneously for all the values in a row of the grid, say in row  $\#k$ .

The line Jacobi method is defined by

$$u_h^{(m+1)}(x_j, y_k) = \frac{1}{4} \left\{ u_h^{(m+1)}(x_{j+1}, y_k) \right. \\ \left. + u_h^{(m+1)}(x_{j-1}, y_k) + u_h^{(m)}(x_j, y_{k+1}) \right. \\ \left. + u_h^{(m)}(x_j, y_{k-1}) \right\} - \frac{h^2}{4} g(x_j, y_k), \quad 1 \leq j \leq N - 1$$

for  $k = 1, \dots, N - 1$ . This is a tridiagonal system of order  $N - 1$ ; and it can be solved in approximately  $5N$  arithmetic operations.

## MULTIGRID ITERATION

I just give the general philosophy, as the details are quite complicated.

Imagine having a sequence of linear systems

$$A_N \mathbf{u}_N = \mathbf{b}_N, \quad N = N_0, N_1, N_2, \dots, N_q$$

In our case with the Poisson equation, we would generally have  $N_{j+1} = 4N_j$ , approximately. We want to solve for the most accurate discretization, that for  $N = N_q$ .

The general philosophy is to use information from each of the systems  $A_N \mathbf{u}_N = \mathbf{b}_N$  to solve all the systems with a finer mesh. There are a number of ways in which this may be carried out.

Step 1. Use the Jacobi method to iterate in  $A_{N_q} \mathbf{u}_{N_q} = \mathbf{b}_{N_q}$ , calculating a certain number of iterates (say  $K$ ), to get  $\mathbf{v} = \mathbf{u}_{N_q}^{(K)}$ . Then calculate the residual

$$\mathbf{R} = \mathbf{b}_{N_q} - A_{N_q} \mathbf{v} = A_{N_q} (\mathbf{u}_{N_q} - \mathbf{v})$$

Step 2. Find an approximate solution by performing the multigrid iteration on the system

$$A_{N_{q-1}} \delta = \mathcal{R}_{q,q-1} \mathbf{R}$$

calling the approximate solution  $\tilde{\delta}$ .

Step 3. *Prolong* the solution  $\tilde{\delta}$  to a vector  $\hat{\delta}$  be of length  $N_q$  and then define the new iterate for  $A_{N_q} \mathbf{u}_{N_q} = \mathbf{b}_{N_q}$  to be

$$\mathbf{u}_{N_q} \approx \mathbf{v} + \hat{\delta}$$

This converges extremely rapidly; and there are public domain packages available which implement it for common classes of partial differential equations. In general these are the fastest means of solving discretizations of PDEs.

## DIRECT METHODS

With the systems for some special PDEs, including the Poisson equation on a rectangle, it is possible to give an exact method of solution which is also extremely efficient. It depends on knowing the eigenvalues and eigenvectors for the matrix  $A_N$ . Since the matrix is symmetric and positive definite, there is a basis of orthogonal eigenvectors; and it can be found explicitly. Suppose it is written as

$$\mathbf{v}^{(1)}, \mathbf{v}^{(2)}, \dots, \mathbf{v}^{(N)}$$

with eigenvalues  $\lambda_1, \dots, \lambda_N$ . Then write

$$\mathbf{b} = \sum_{j=1}^N (\mathbf{b}, \mathbf{v}^{(j)}) \mathbf{v}^{(j)}$$

The solution of  $A\mathbf{u}_N = \mathbf{b}$  is given by

$$\mathbf{u}_N = \sum_{j=1}^N \frac{1}{\lambda_j} (\mathbf{b}, \mathbf{v}^{(j)}) \mathbf{v}^{(j)}$$

In the particular case of the Poisson equation over a rectangle, this can be computed in  $O(N \log N)$  operations.

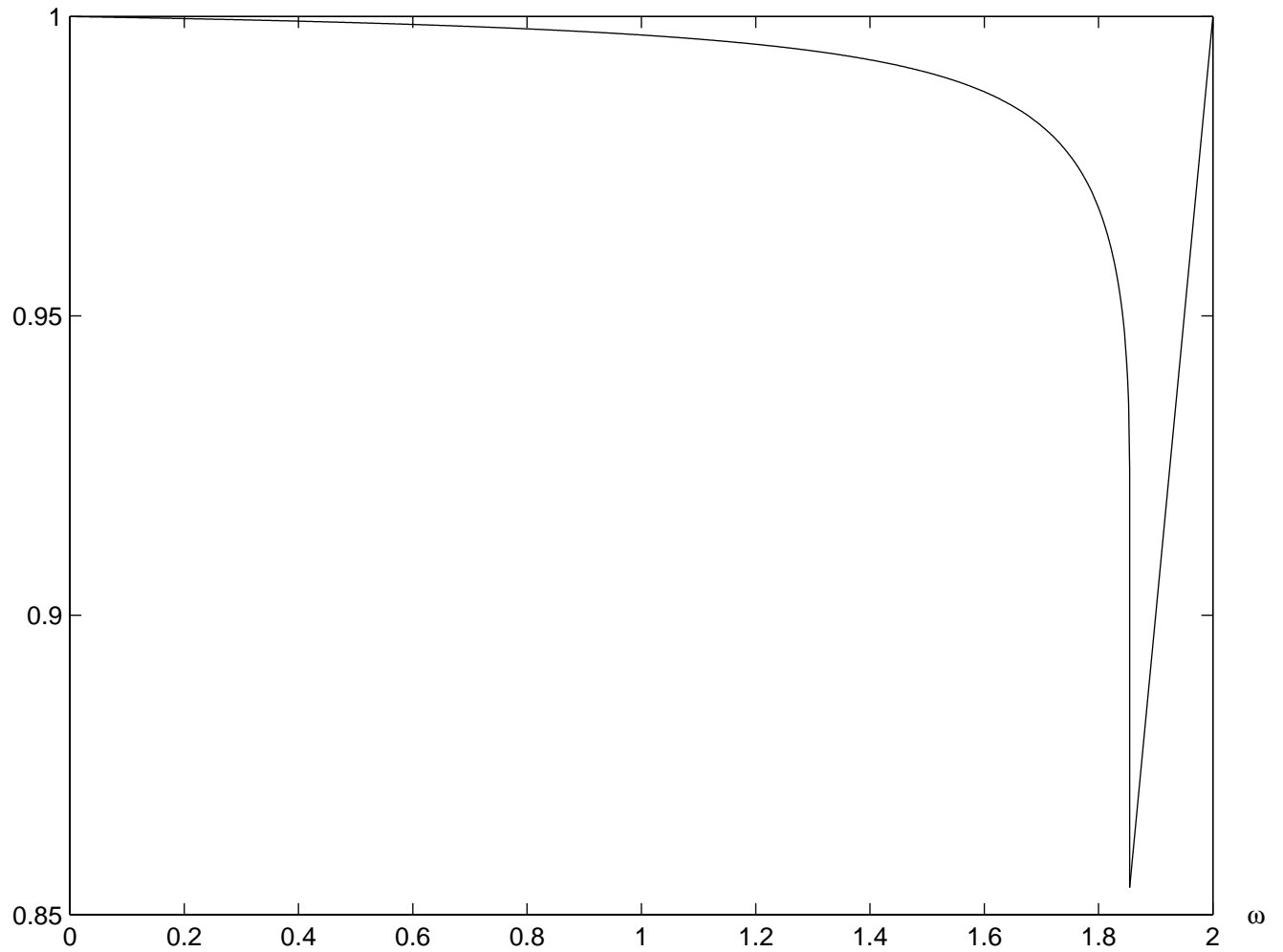


Figure 1: Convergence rate  $r_\sigma(M(\omega))$  with  $N = 40$  and  $\omega^* = 1.85450$

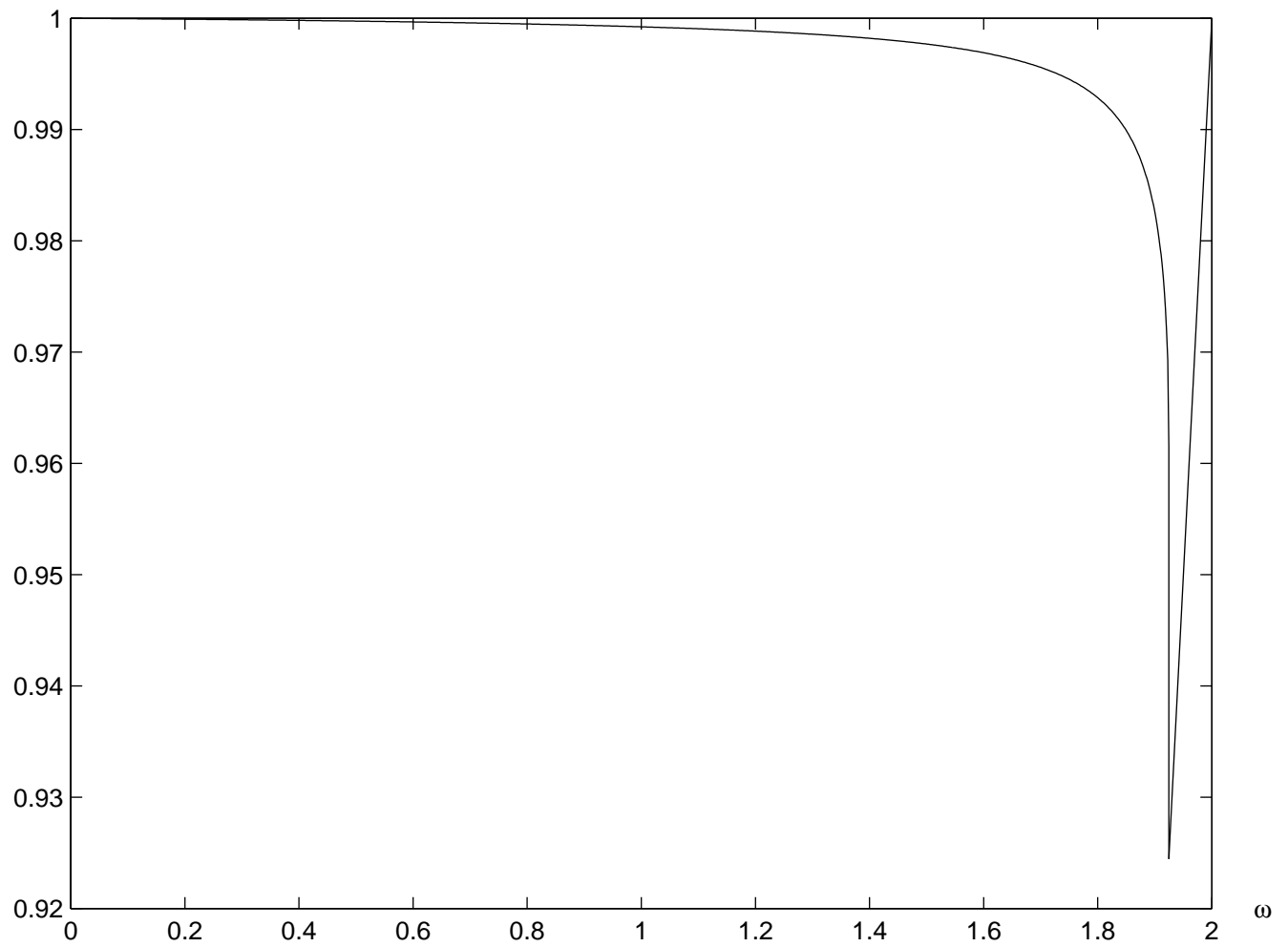


Figure 2: Convergence rate  $r_\sigma(M(\omega))$  with  $N = 80$  and  $\omega^* = 1.92445$